



EVOLUTION IN THE SYSTEM CORONA-DISK

Krasimira Dimitrova Yankova
 Space Research and Technology Institute – BAS
 Space Astrophysics Department
 E-mail: f7@space.bas.bg

X SBAC
 30 May – 3 June, 2016,
 Belgrade, Serbia

ABSTRACT

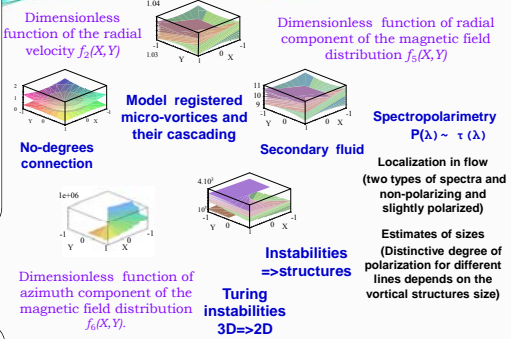
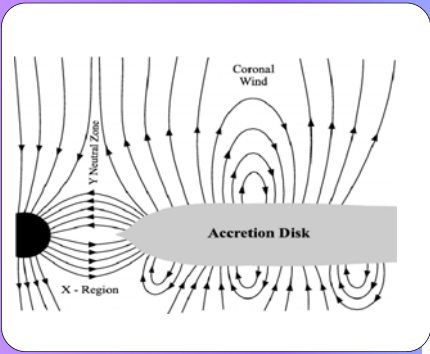
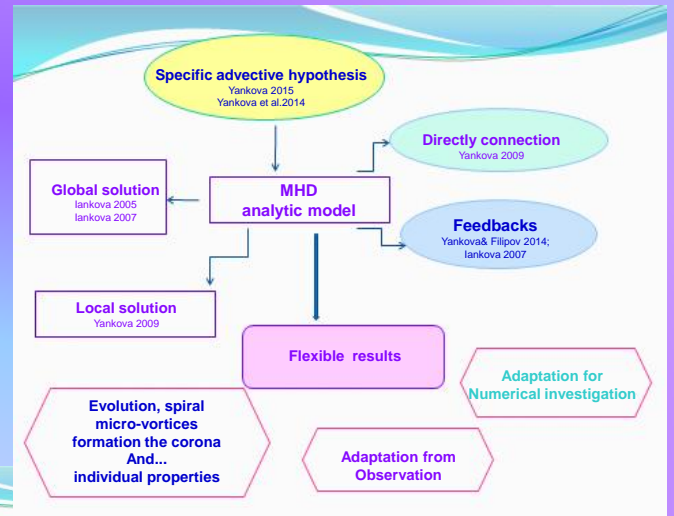
This paper considers magneto-hydrodynamics of the system disk-corona. We will researching the structuring of the flow on the secondary component. Analyze the influence of the distributions on boundary with the primary component. Discusses the significance of the type of the border for exchange of energy between components.

Unified model of the AGN
 A. A. Abdo et al. 2009

<http://www.uni-goettingen.de/en/216897.html>



Group of 140 (one hundred and forty) authors started an observation program whose main conclusions are that no matter the host galaxy a quasar core shows similar structure and the same mechanism of development and differences in the observations are the result of different levels of accretion, mass and direction monitoring.



Boundary corona-disk is especially important because there manifest themselves the effects of the warming into the pad: Tightening of advective rings to the center in disk and negative entropy realized a new state. For radiative corona the warming is a major supports factor and does not allow corona to attenuation.

These results are another step in the direction of for creating full model of the system disk-corona. Also, tool for determining the degree to which the main component influenced the development of coronary component.

Basic model equations of the magneto-hydrodynamics of accretion - disk flow

$$\begin{aligned}
 \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) &= 0 & \nabla \cdot \mathbf{v} &= 0 & \nabla \cdot \mathbf{B} &= 0 \\
 \frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} &= -\frac{1}{\rho} \nabla p - \nabla \Phi + \left(\frac{\mathbf{B}}{4\pi\rho} \cdot \nabla \right) \mathbf{B} + \mathcal{G} \nabla^2 \mathbf{v} \\
 \frac{\partial \mathbf{B}}{\partial t} &= \nabla \times (\mathbf{v} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B} & \Phi &= -\frac{GM}{r - r_g} \\
 \rho \Gamma \frac{\partial S}{\partial t} - \frac{M}{2\pi r} \Gamma \frac{\partial S}{\partial r} &= Q^+ - Q^- + Q_m & \eta &= \frac{\eta_m}{\rho} = \frac{c^2}{4\pi\sigma} \\
 p &= p_r + p_\theta + p_m & r_g &= \frac{2GM}{c^2}
 \end{aligned}$$

$$v_r \frac{\partial \rho}{\partial r} + v_z \frac{\partial \rho}{\partial z} + \rho(\omega + k_\phi \Omega) = 0$$

$$\frac{\partial B_r}{\partial r} = \mu \frac{5r^2 + z^2}{(r^2 + z^2)^2} \sqrt{4r^2 + z^2} - \frac{B_r + k_\phi B_\phi}{r}$$

$$v_r \frac{\partial v_r}{\partial r} + v_z \frac{\partial v_r}{\partial z} + v_r(\omega + k_\phi \Omega) = \frac{\partial v_r^2}{\partial r} - \frac{v_r^2}{\rho} \frac{\partial \rho}{\partial r} + \frac{k_\phi B_r B_\phi}{4\pi\rho r} + \frac{B_r}{4\pi\rho} \frac{\partial B_r}{\partial z}$$

$$v_r \frac{\partial}{\partial r} (\Omega r^2) + v_z \frac{\partial}{\partial z} (\Omega r^2) + (\Omega r^2) \omega = \frac{B_r}{4\pi\rho r} \frac{\partial}{\partial r} (r^2 B_\theta) + \frac{B_r r}{4\pi\rho} \frac{\partial B_\theta}{\partial z}$$

$$v_r \frac{\partial v_r}{\partial r} + v_z \frac{\partial v_r}{\partial z} + v_r(\omega + k_\phi \Omega) = \frac{\partial v_r^2}{\partial r} - \frac{v_r^2}{\rho} \frac{\partial \rho}{\partial r} + \frac{\mu B_r}{2\pi\rho} \frac{z^4 + 5z^2 r^2 - 2r^4}{(r^2 + z^2)^2} \sqrt{4r^2 + z^2}$$

$$\frac{\partial v_r}{\partial r} + \frac{v_r}{r} + \frac{\partial v_z}{\partial z} = 0 \quad \frac{aT^4}{3\rho} = v_r^2 - v_z^2 / 2 - RT$$

Model equations of the fluid in the disk's corona, for sharp boundary between components

$$\begin{aligned}
 f_1(x, H) &= -\frac{c_1 + c_2}{(1 - c_{12} - c_{14})^2} \left[\frac{1}{x^2} - \frac{1}{(x-x_1)^2} \right] + 1 \\
 f_2(x, H) &= (9c_{10} + 2)(1 - c_{12} - c_{14})(x-1) + 1 \\
 f_3(x, H) &= \frac{c_2}{4} (1 - c_{12} - c_{14})^2 (\alpha^2 - 1) x^2 + \frac{c_{10}}{2} (1 - c_{12} - c_{14})^2 (\alpha^2 - 1) x^2 + \\
 &+ \frac{c_1}{4} (1 - c_{12} - c_{14})^2 (\alpha^2 - 1) + (c_1 + c_2) \left(\frac{1 - c_{12} - c_{14}}{2(1 - c_{12} - c_{14})^2} \right) (x - x_1 - 1) + 1 \\
 f_4(x, H) &= \frac{1 + c_2}{3} (1 - c_{12} - c_{14}) \left(\frac{1}{x^2} - 1 \right) + 1 \\
 f_5(x, H) &= (1 - c_{12} - c_{14})^2 \left[\frac{3 + \frac{2c_{10}}{4c_1 x^2} (x - x_1 - 1) - \frac{2c_{10}}{c_1 x^2} \frac{1}{(x-x_1)}}{1 + \frac{c_{10}}{c_1 x^2} \frac{1}{(x-x_1)} + \frac{c_{10}}{c_1 x^2} \frac{1}{(x-x_1)^2}} \right] \\
 f_6(x, H) &= \frac{c_1(1 - c_{12} - c_{14})^2}{c_1} \left(\frac{1}{x^2(x-x_1)} - \frac{1}{x^2} \right) + \frac{c_{10}(1 - c_{12} - c_{14})^2}{c_1} \left(\frac{1}{x^2(x-x_1)} - \frac{1}{x^2} \right) + \\
 &+ \frac{5 + 20c_{10}(1 - c_{12} - c_{14})}{c_1} \left(\frac{1}{x} - 1 \right) - \frac{4c_{10}}{c_1} \left(\frac{1}{x} - 1 \right) \\
 f_7(x, H) &= \left(\frac{c_1 + 1}{k_\phi} + \frac{2c_2}{c_{14}} \right) (x-1) + \frac{(c_1 + c_2)}{c_1(1 - c_{12} - c_{14})^2} (x-x_1) - \\
 &+ \frac{k_\phi + 2}{2k_\phi(1 - c_{12} - c_{14})^2} \left(\frac{1}{x^2} \right) + \frac{c_2}{c_1(1 - c_{12} - c_{14})^2} \left(\frac{1}{x^2} \right) + \\
 &+ \frac{1 + c_2}{(1 - c_{12} - c_{14})} \left(\frac{1}{x} \right) \\
 f_8(x, H) &= \frac{c_2}{(x-x_1)} - \frac{c_2(1-x^2)}{2(1 - c_{12} - c_{14})^2} - \frac{c_1(1 - c_{12} - c_{14})^2}{3} (1 - c_{12} - c_{14}) + 1
 \end{aligned}$$

Boundary distributions from results of 3D-model of the disk, for sharp boundary