

Interpreting spectropolarimetric results across a broad spectral range - 3D radiative transfer with the STOKES code

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Talk at the 9th Serbian Conference
on Spectroscopic Line Shapes in Astrophysics (SCSLSA)

16th May 2013

Banja Koviljača, Serbia



Credits to our polarization network



Astronomical Institute
Academy of Sciences
of the Czech Republic



Vladimír Karas
Michal Dovčiak



Fabio Muleri



Giorgio Matt



C. Martin Gaskell



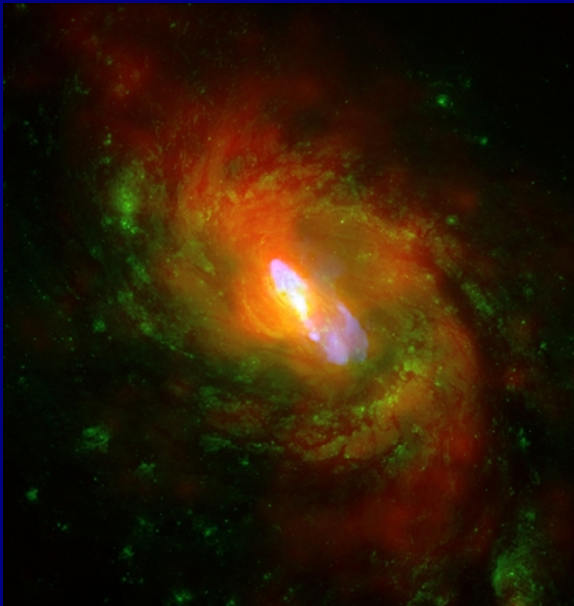
Delphine Porquet
Frédéric Marin
Francesco Tamborra
Nicolas Grosso
René W. Goosmann

Why should we care about polarization?

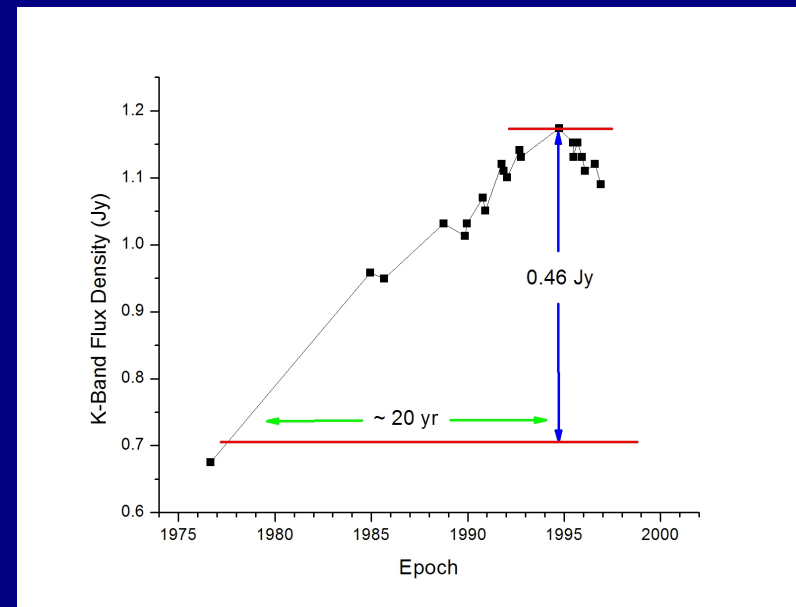
We practice observational astronomy *mainly* based on electromagnetic (EM) radiation.

The EM radiation tells us about its emission processes and its interactions with matter.

The information is usually exploited as a function of wavelength, time, and space → **(time-resolved) spectroscopy and imaging**



X-ray (NASA/CXC/MIT/C.Canizares, D.Evans et al), Optical (NASA/STScI), Radio (NSF/NRAO/VLA)

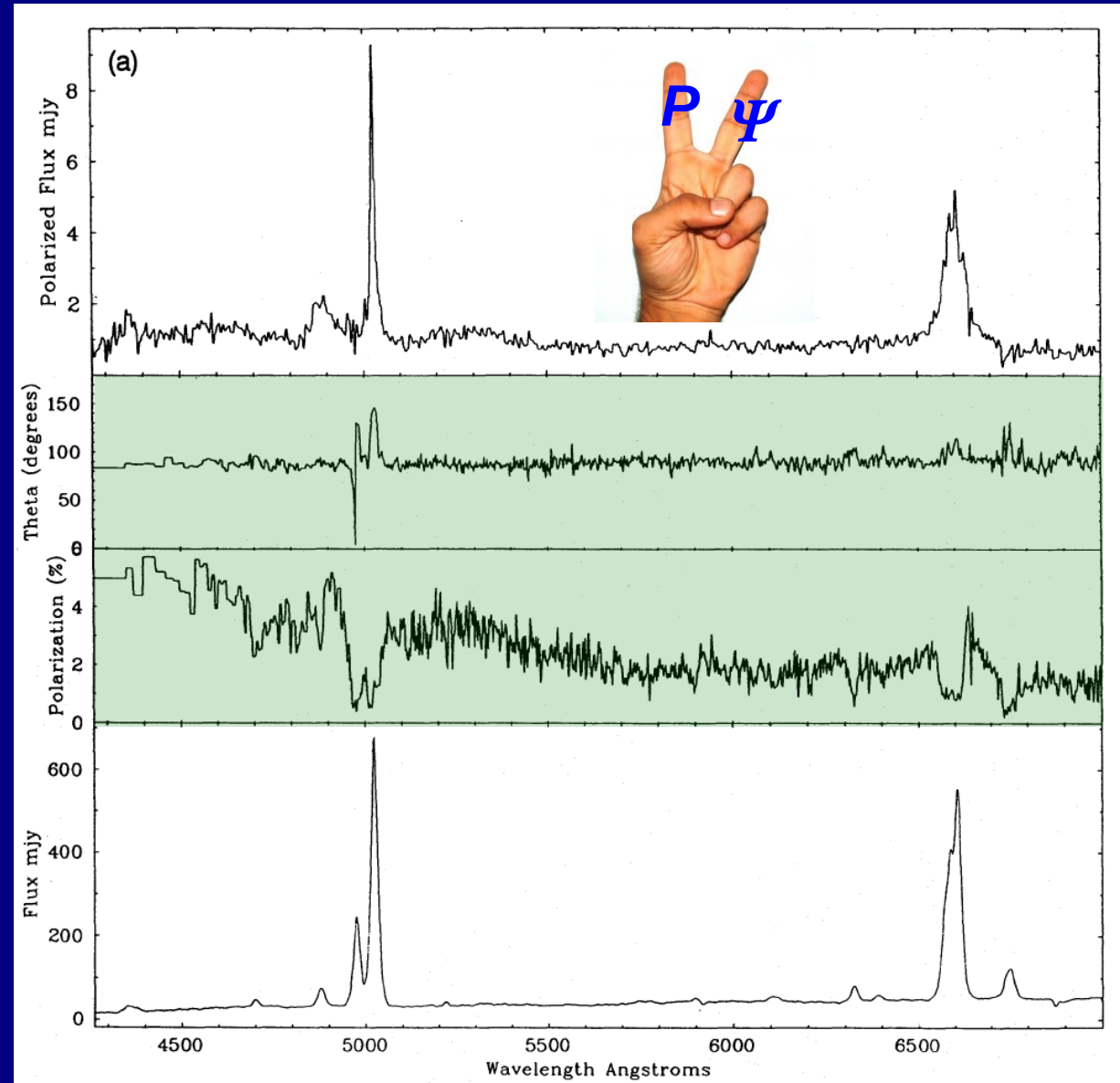


2.2 micron light curve of NGC 1068 (work by I. Glass)

Why should we care about polarization?

BUT: almost any interaction of EM radiation with matter also modifies its polarization state!

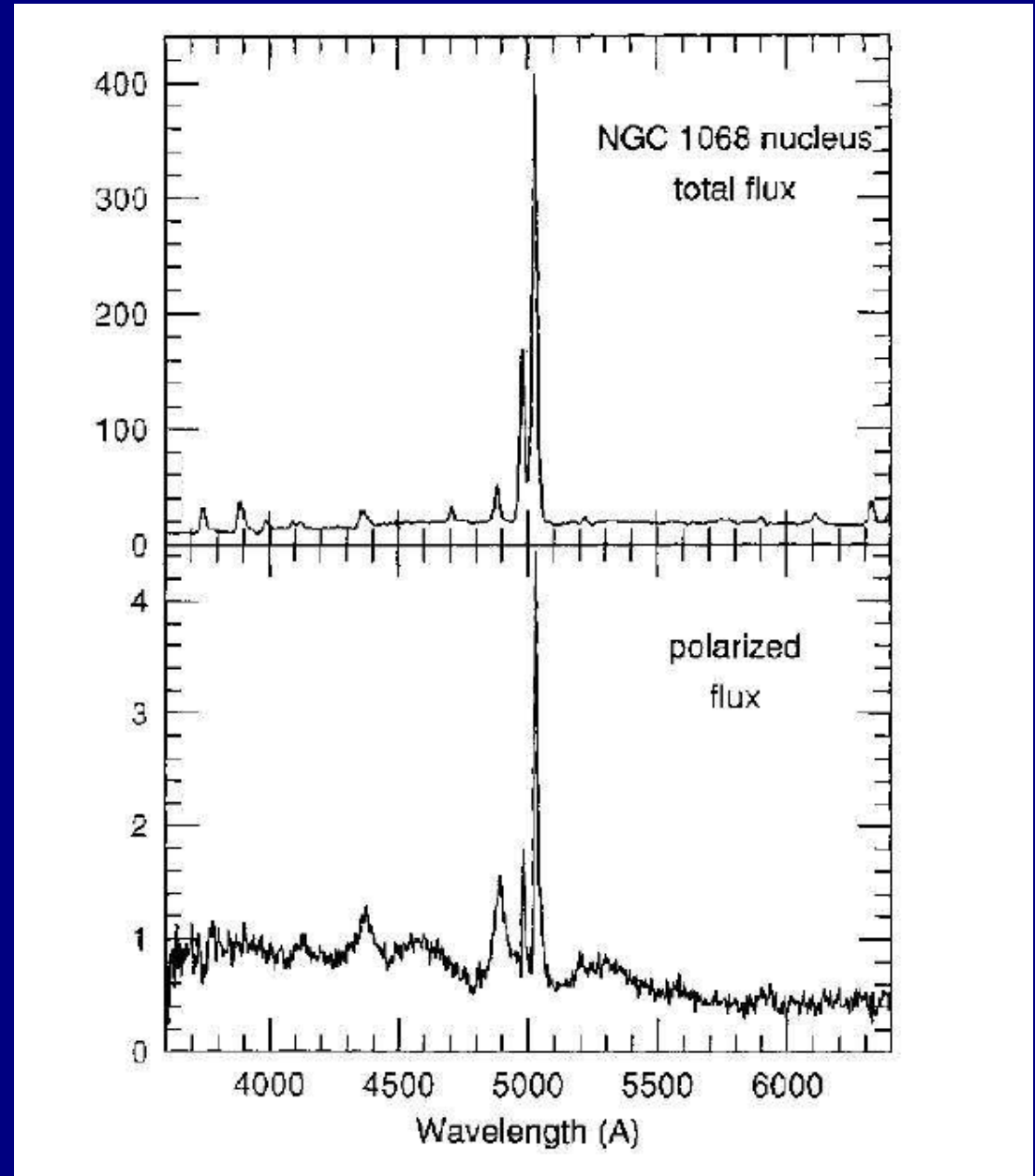
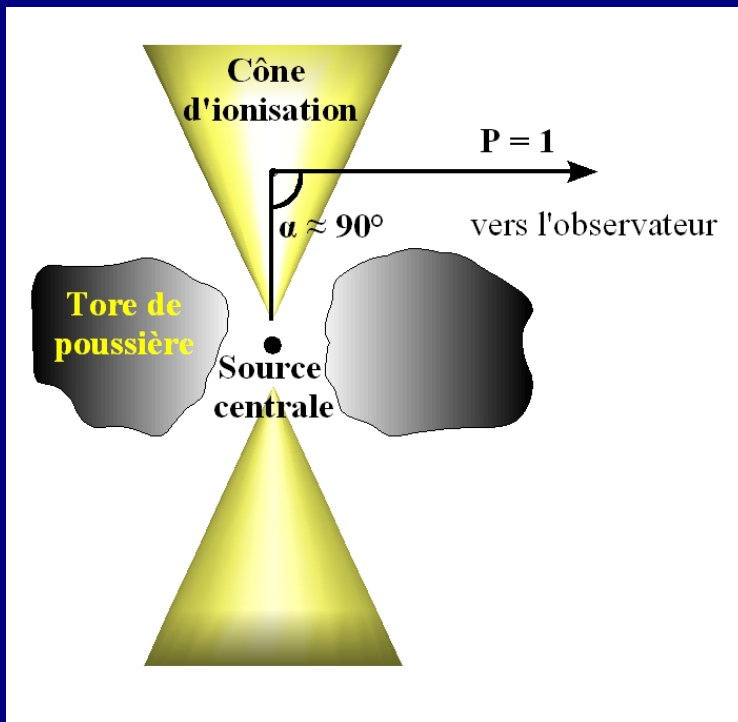
ERGO: Considering the polarization state of light gives us a set of **two additional, independent observables** as a function of photon wavelength, time, and space.



Radio-quiet objects Hidden type-1 AGN

A major break-through for the unified model for NGC 1068 (Antonucci & Miller 1985)

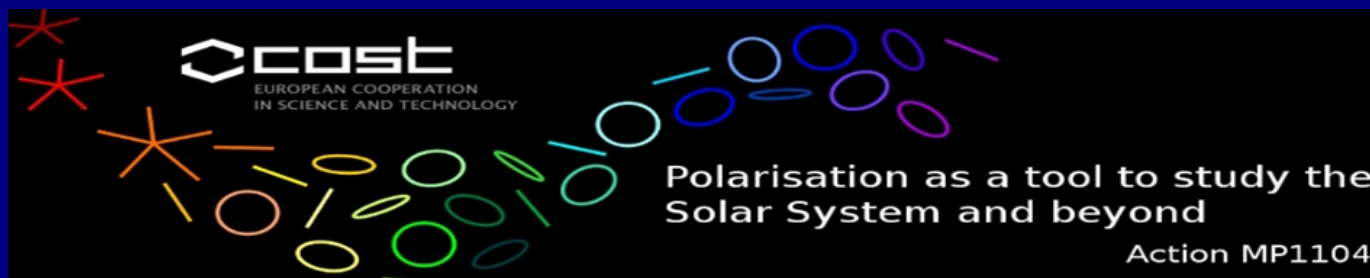
→ periscope view of AGN in polarized flux



Exploiting and modeling polarization in astrophysics

- The basic concept of (linear) polarization
- Techniques to observe broad-band polarization
- Mechanisms producing polarization
- Modeling scattering-induced polarization with *STOKES*
- Summary and conclusions

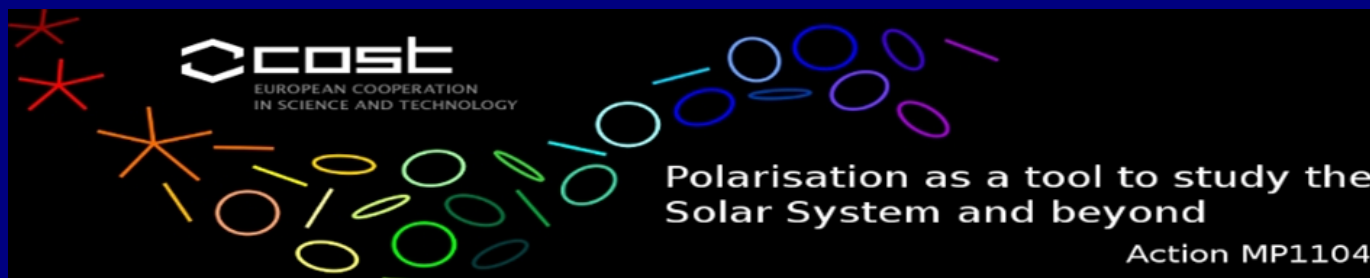
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The root of all electromagnetism...

Maxwell's equations

$$\nabla \cdot \mathbf{E} = \rho / \epsilon_0$$

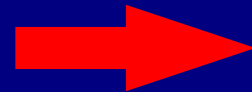
$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{j}_c$$

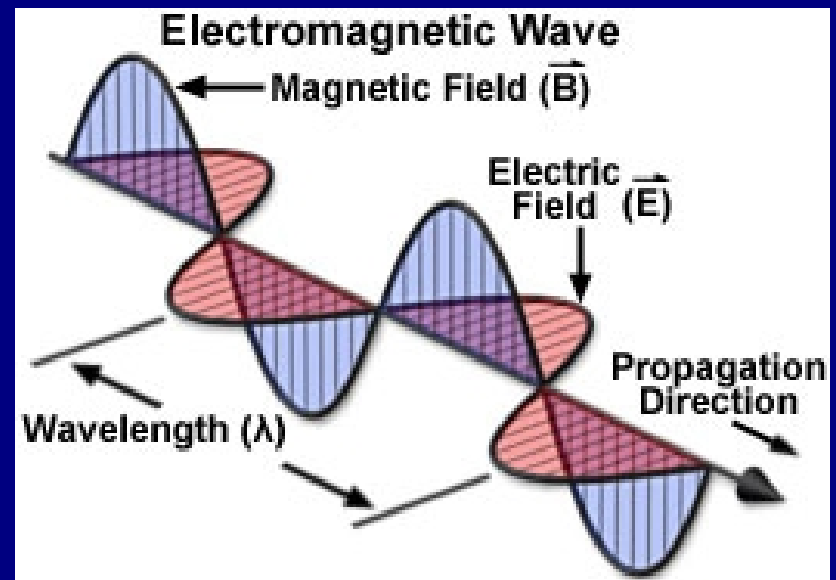
where

$$\nabla = \hat{\mathbf{i}} \frac{\partial}{\partial x} + \hat{\mathbf{j}} \frac{\partial}{\partial y} + \hat{\mathbf{k}} \frac{\partial}{\partial z}$$



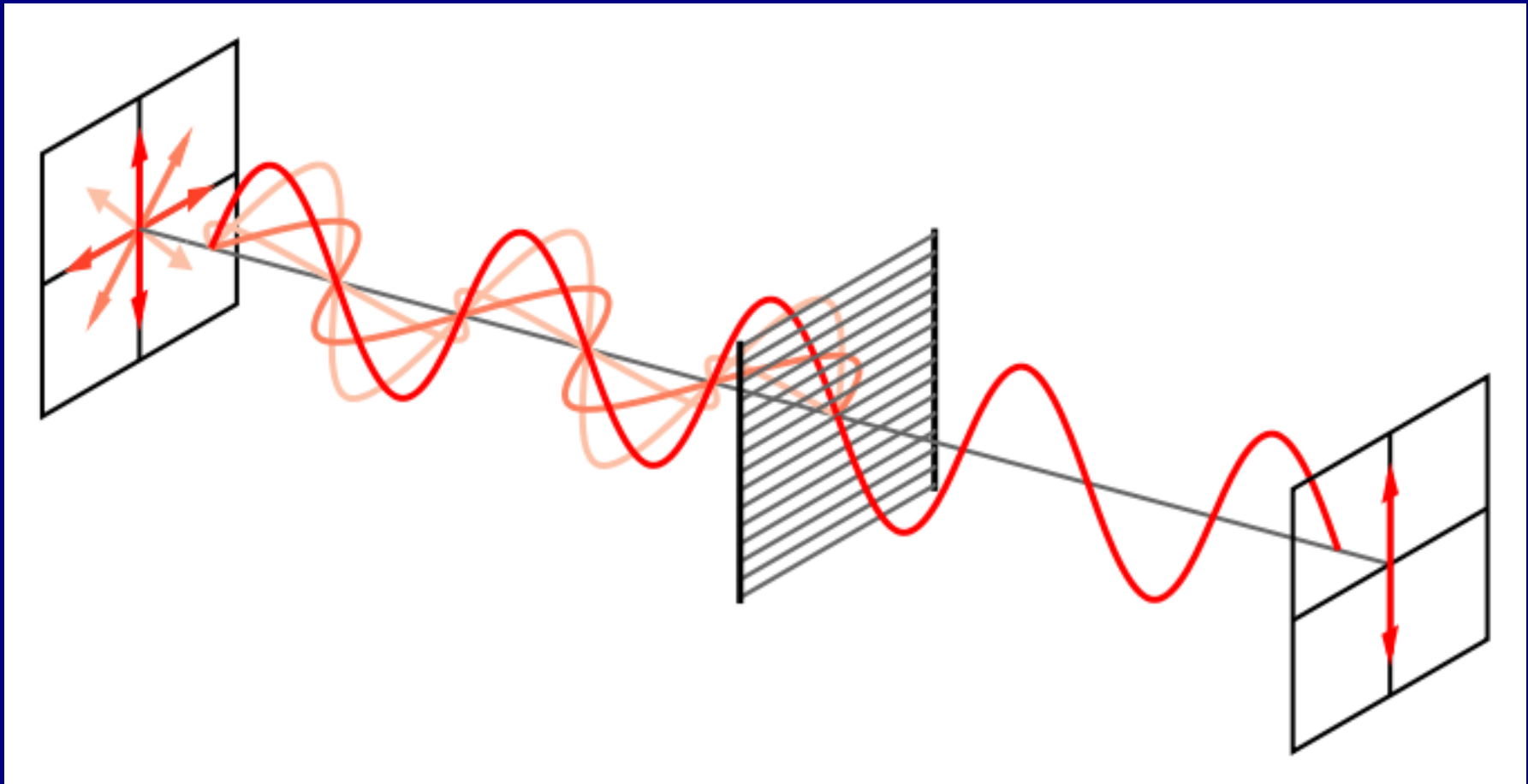
3D wave equation for the electric field

$$\nabla^2 \mathbf{E} = \mu_0 \epsilon_0 \frac{\partial^2 \mathbf{E}}{\partial t^2}$$



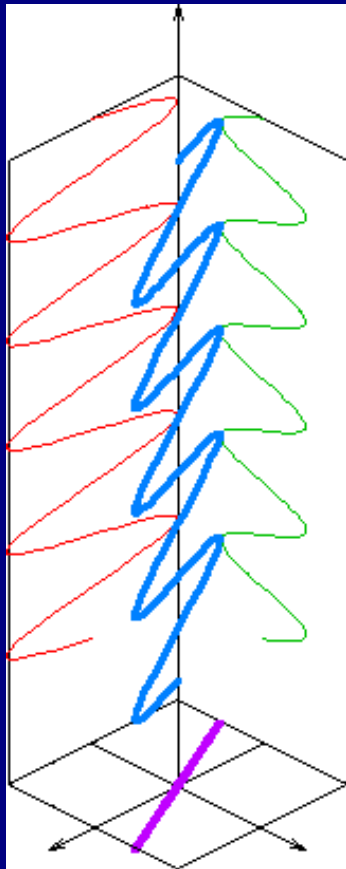
Linear polarization of light

The polarization state of an electromagnetic wave denotes the direction of the electric field vector (classical picture)

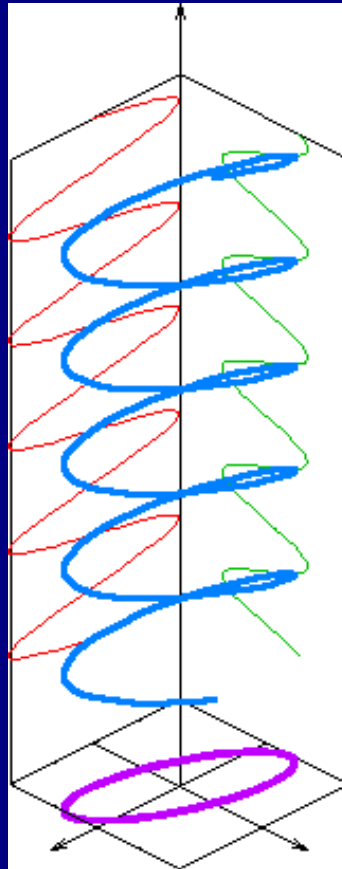


Polarization of coherent light

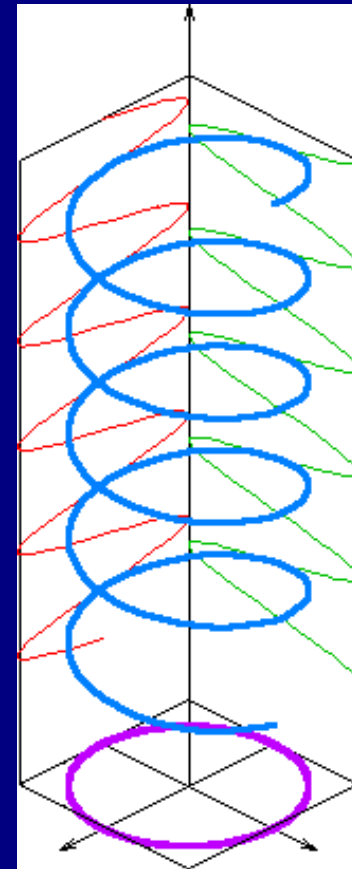
A **coherent** electromagnetic wave can be decomposed in two perpendicular components with a defined phase relation.



linear



elliptical



circular

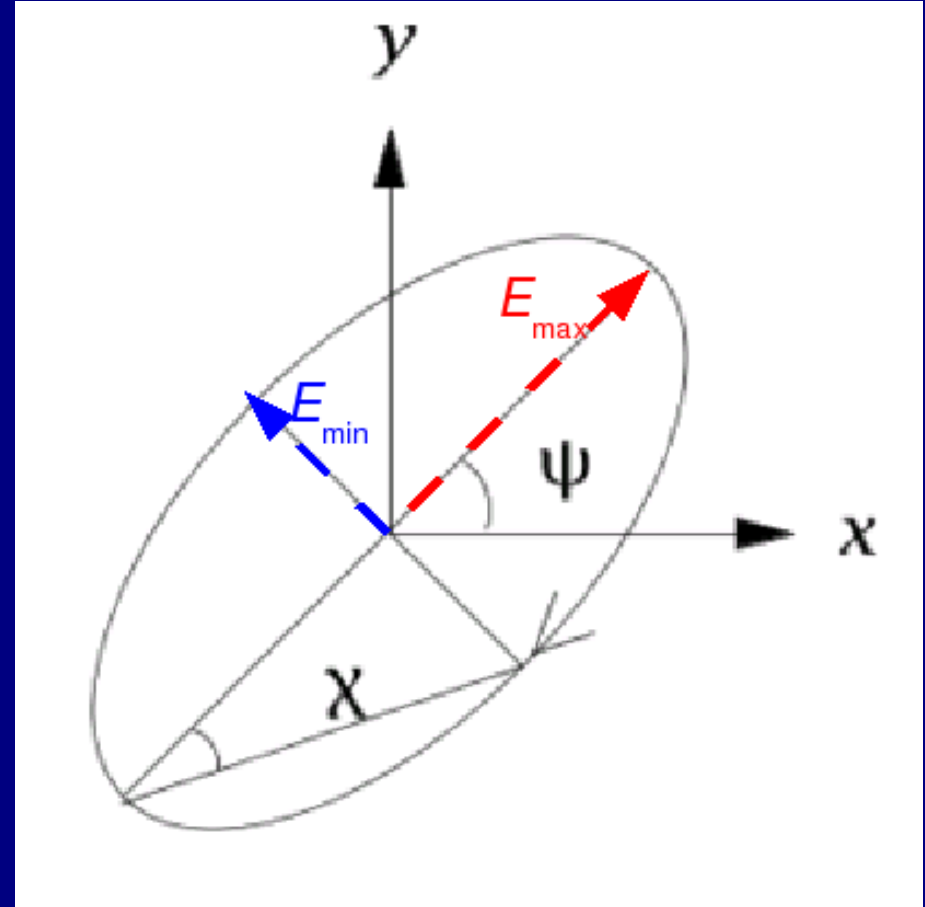


The polarization ellipse

The linear polarization degree P is defined by

$$P = \frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}}.$$

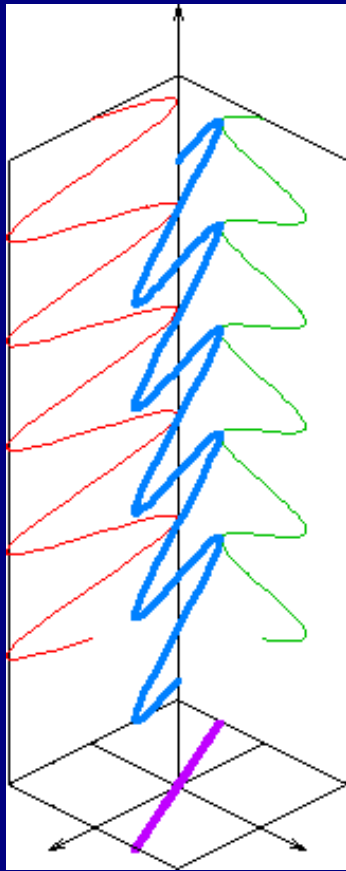
Note: $0 \leq P \leq 1$.



Herein, I_{\max} and I_{\min} are measured along the directions at which the length of the E -vector has a maximum or minimum, respectively.

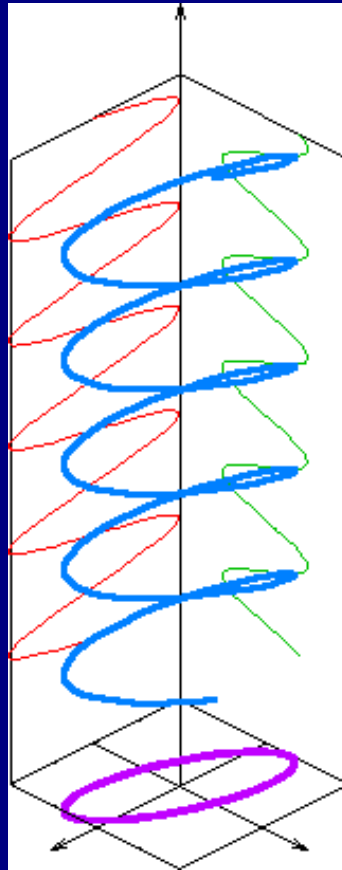
Polarization of coherent light

100%
0%



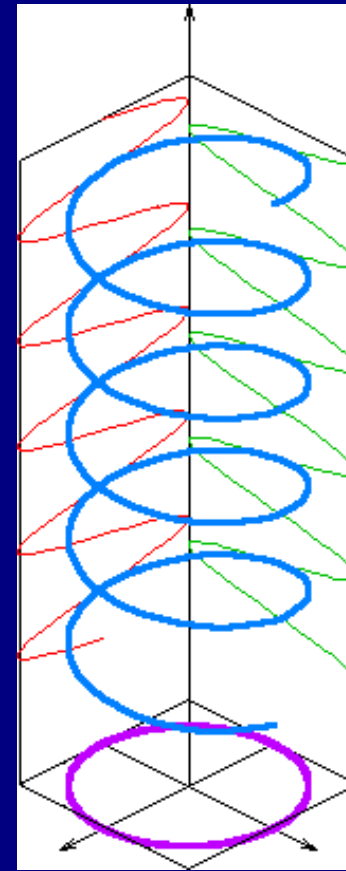
linear

~30%
~70%



elliptical

0%
100%



circular

Linear pol.
Circular pol.



Astronomical light and (none-)coherence

The light from astronomical sources comes from uncorrelated sub-sources :

- different parts of a stellar surface,
- different layers inside an ionized nebula,
- different distances in redshift
- ...

There is only a statistical coherence of “astronomical light” and a linear polarization degree of 100% is quasi-impossible to observe.

For the remainder of this presentation we only discuss linear polarization that is more important in most wave bands than circular polarization.

The Stokes parameters

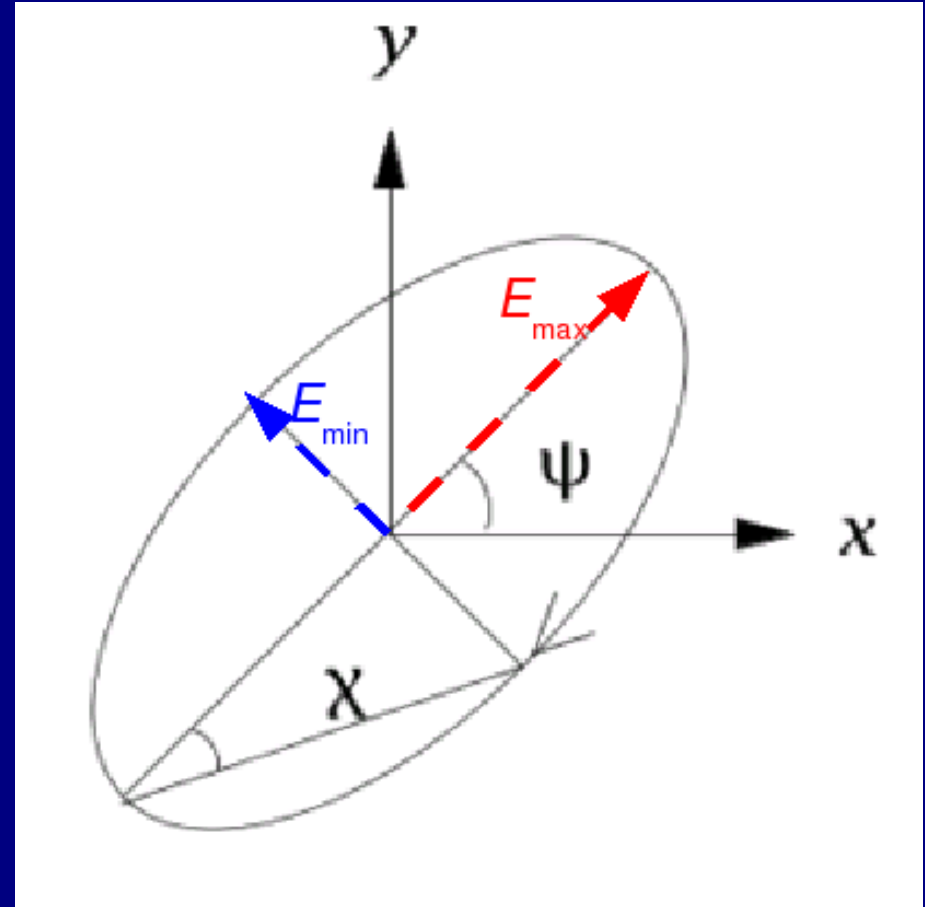
The polarization state is completely described by the Stokes parameters:

$$I = \langle E_{max}^2 + E_{min}^2 \rangle,$$

$$Q = \langle (E_{max}^2 - E_{min}^2) \cos(2\psi) \rangle,$$

$$U = \langle (E_{max}^2 - E_{min}^2) \sin(2\psi) \rangle,$$

$$[V = \langle 2 E_{max} E_{min} \rangle].$$

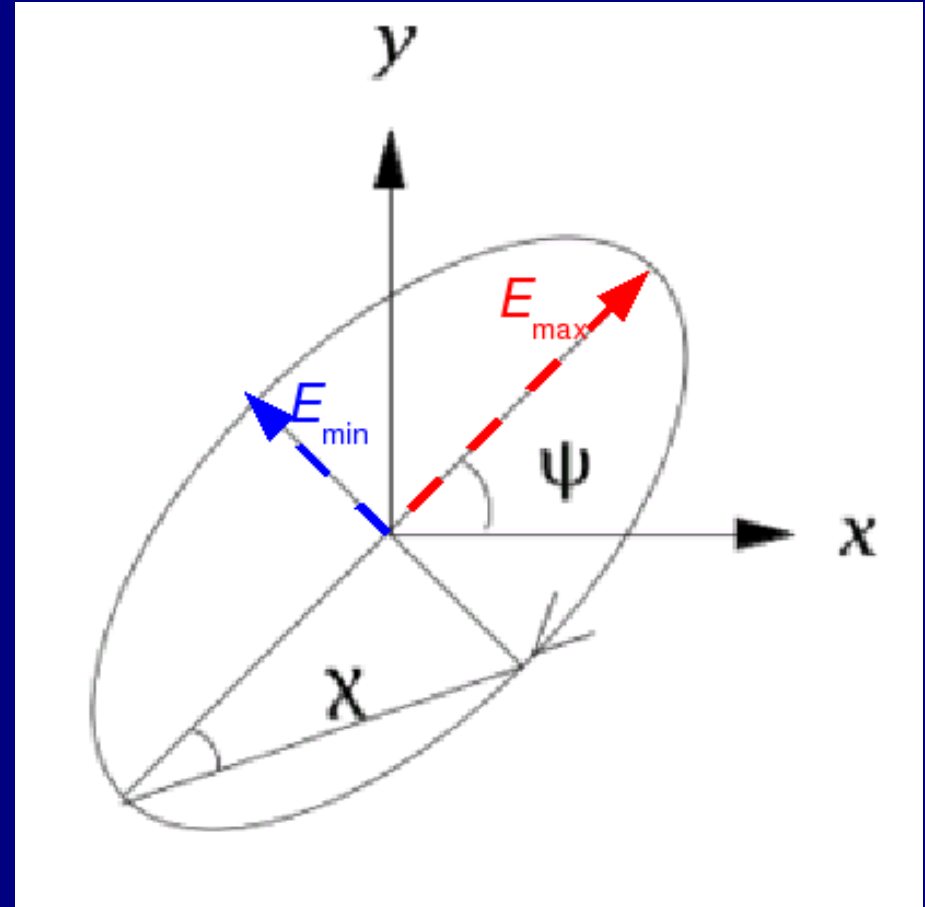


The Stokes parameters

From the Stokes parameters the linear polarization degree P and position angle can Ψ easily be recovered:

$$P = \frac{\sqrt{Q^2 + U^2}}{I},$$

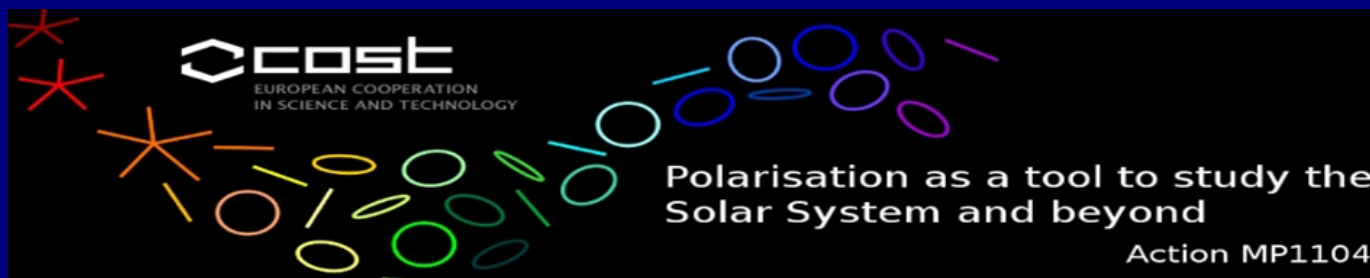
$$\psi = \frac{1}{2} \arctan \frac{U}{Q}.$$



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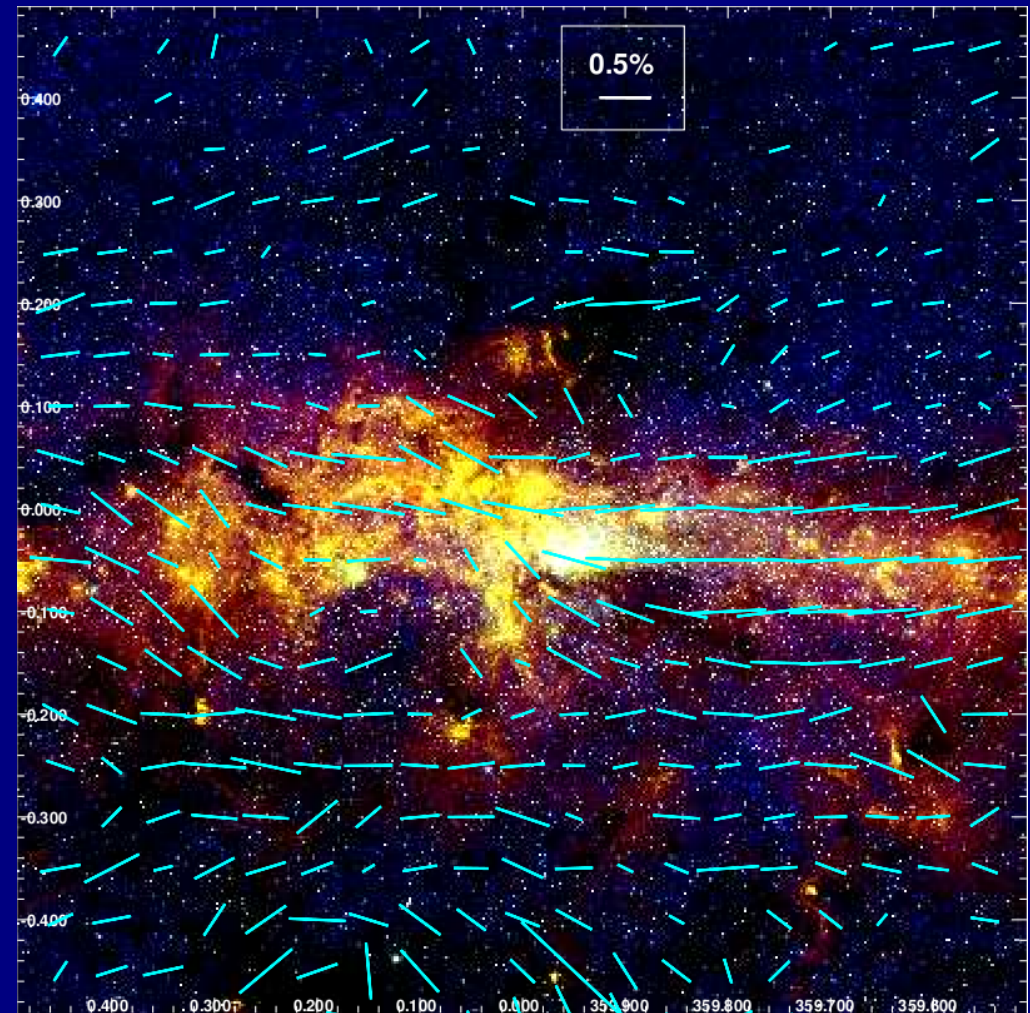


Observing multi-wavelength polarization



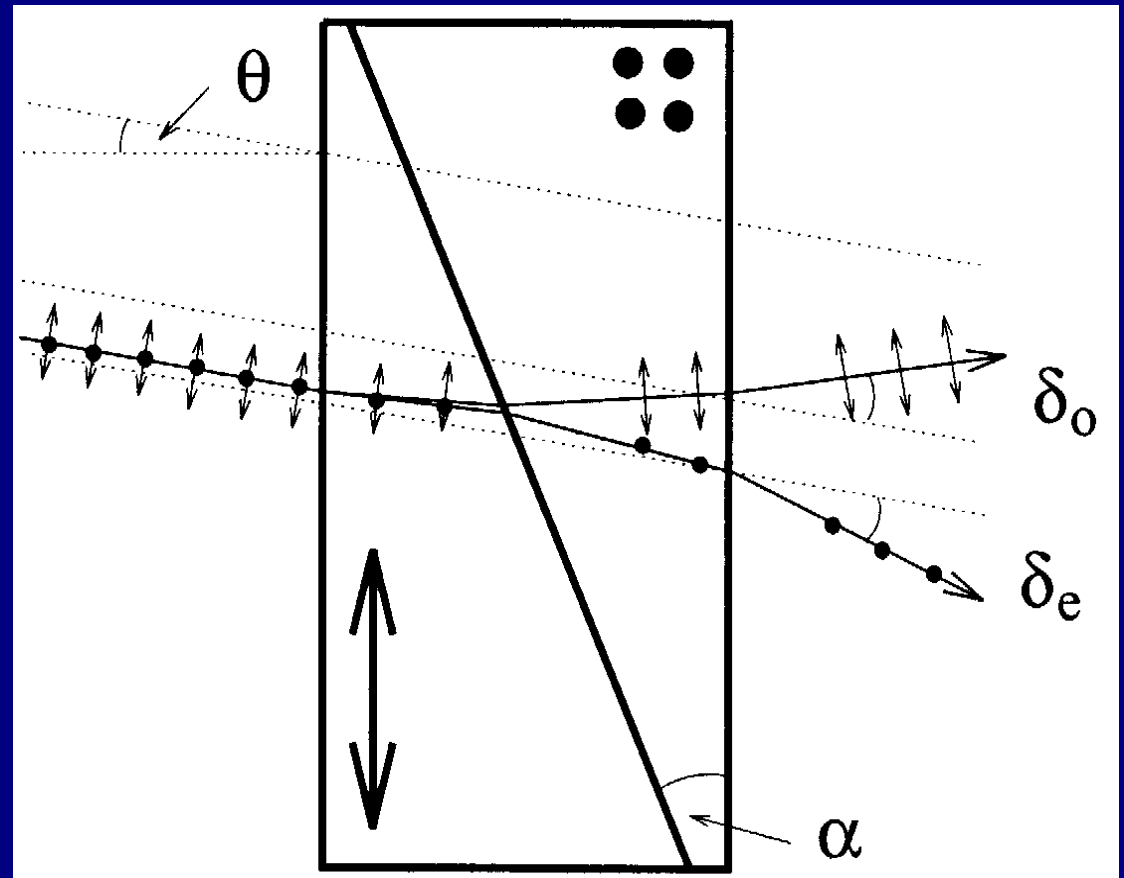
In the radio band, polarimetry comes automatically and for free (position of the antenna)

Infrared polarimetry using rotating wave-plates



Observing multi-wavelength polarization

- In the **optical wave band** It is key to **separate the two perpendicular polarization components** of the light beam.
- This is done using two blocks of **birefringent** material.
- In a birefringent block the speed of light and the refraction properties depend on the polarization orientation.



Wollaston prism (Olivia et al. 1997)

Observing multi-wavelength polarization



X-ray polarimetry relies on micro gas pixel detectors (and future X-ray missions...)

In the optical/UV, birefringent beam-splitters are used to separate the ordinary from the extraordinary beam



Polarimetry capabilities at the VLT



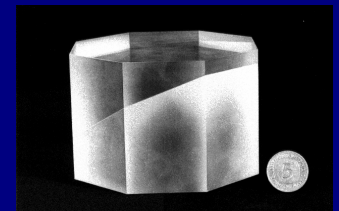
ISAAC mounted on UT3
J, H, K polarimetry



FORS2 mounted
on the VLT
310 nm – 1100 nm
spectropolarimetry
Wollaston prism



NACO mounted on UT4
1-5 micron imaging polarimetry
Wollaston prism



Observational prospects – ready-to-fly technology

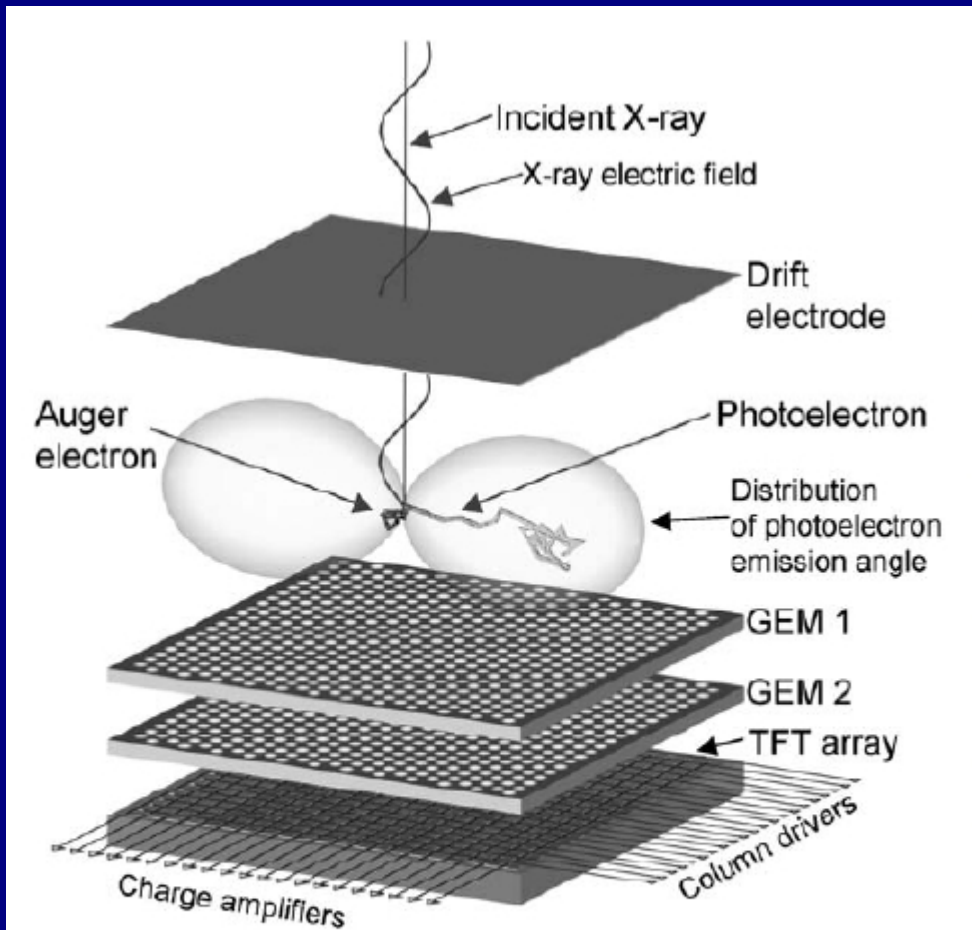


Fig. 1. Schematic diagram of detector geometry used in these measurements. The $\sin^2\theta \cos^2\phi$ distribution of photoelectron emission for normally incident X-rays is projected onto the detector plane and observed as $\cos^2\phi$.

Photoelectric ionization of and subsequent Auger effect
Photo electron and Auger electron both know about the initial polarization of the incident photon

Active-matrix pixel prop. counter

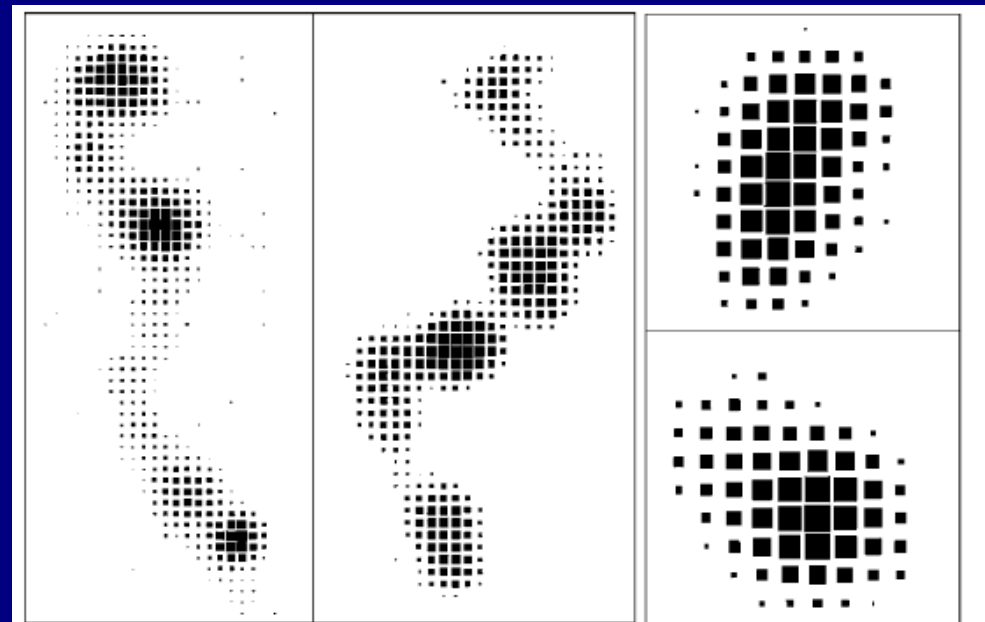
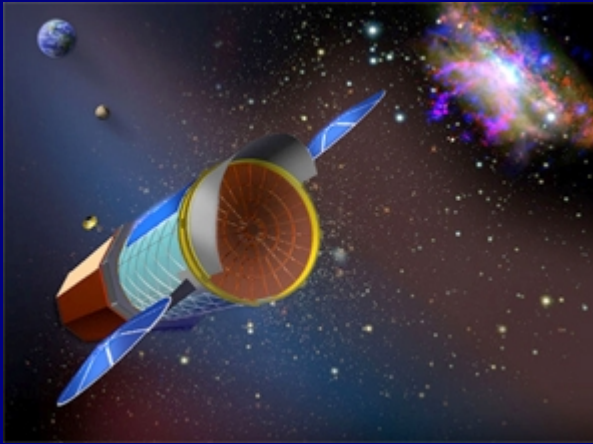


Fig. 3. Track images from 20 keV X-rays (left) and 4.5 keV X-rays (right).

Costa et al. (2001), Bellazzini et al.(2009), Muleri et al. (2009)

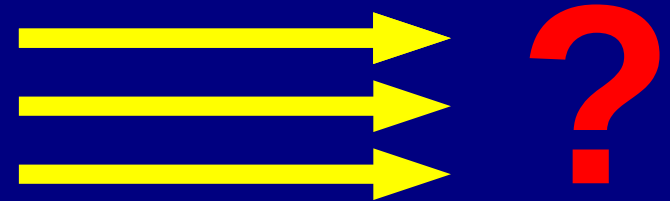
The X-ray polarimetry dilemma : the so-far missing observational perspective



IXO+XPOL (†2011)



GEM SMEX (†2012)



NHXM (†2012)

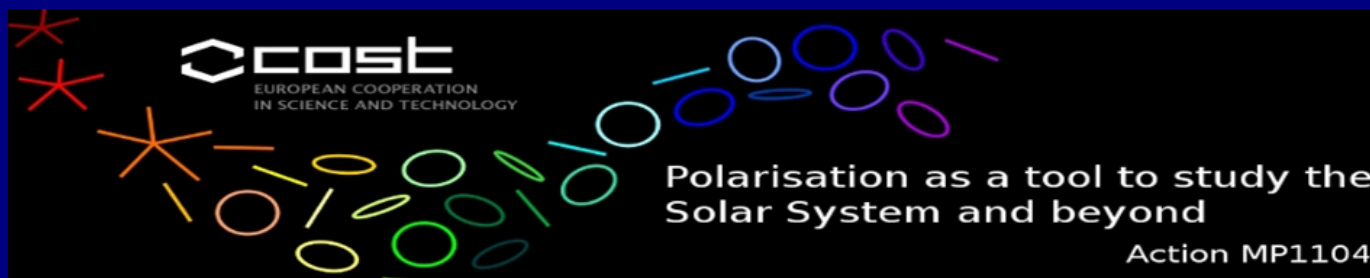


XIPE (†2012)

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Processes producing (de-)polarization

Synchrotron emission

Electron scattering

Dust (Mie) scattering

Resonant line scattering

Dichroic absorption

Faraday rotation

Zeeman line splitting

Dilution (by unpolarized radiation)

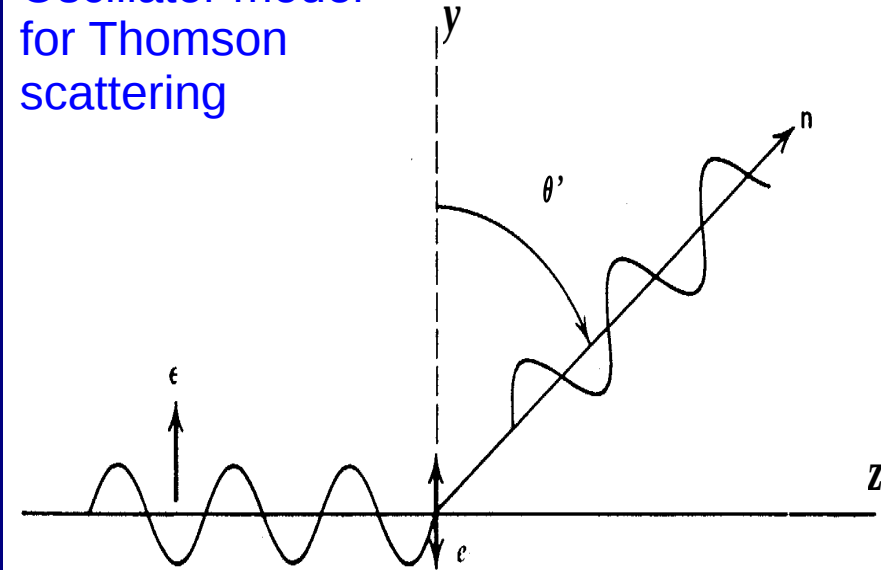
General Relativity

Scattering

Strong polarization: $\Theta = 90^\circ$ (Reflection)

Weak polarization: $\Theta = 0^\circ$ (Transmission)

Oscillator model
for Thomson
scattering



$$\frac{\partial \sigma}{\partial \omega}(\alpha)_{tot} = \frac{1}{2} r_0 (1 + \cos^2 \theta).$$

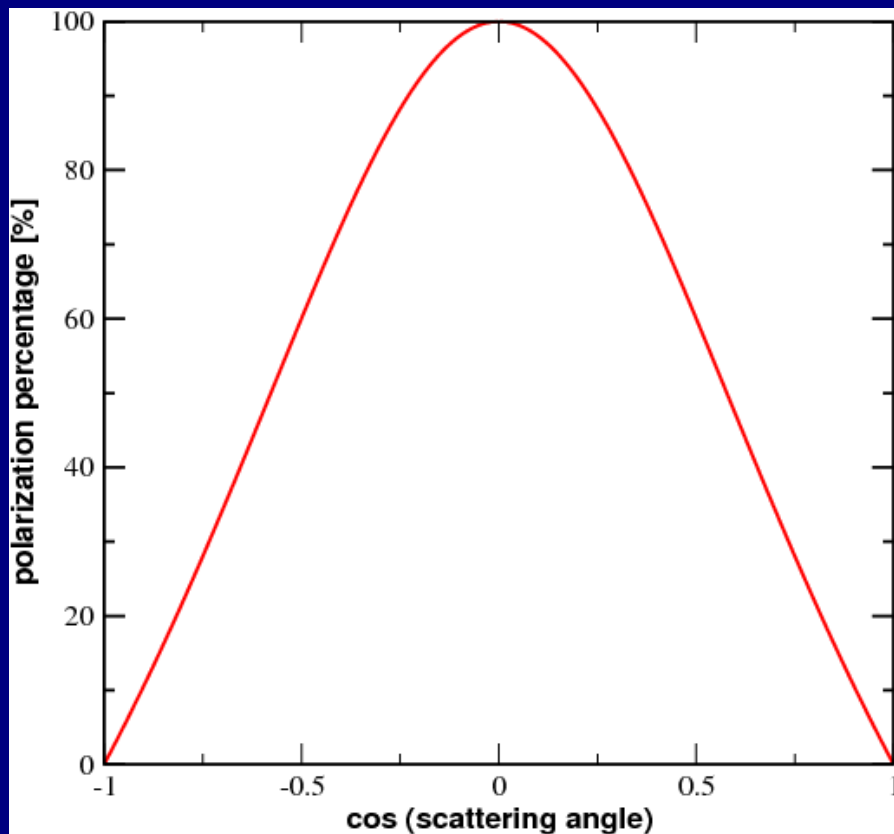
$$P = \frac{1 - \cos^2 \theta}{1 + \cos^2 \theta}.$$

$$\sigma_T = \frac{8\pi}{3} r_0^2 = \frac{8\pi e^4}{3m^2 c^4}.$$

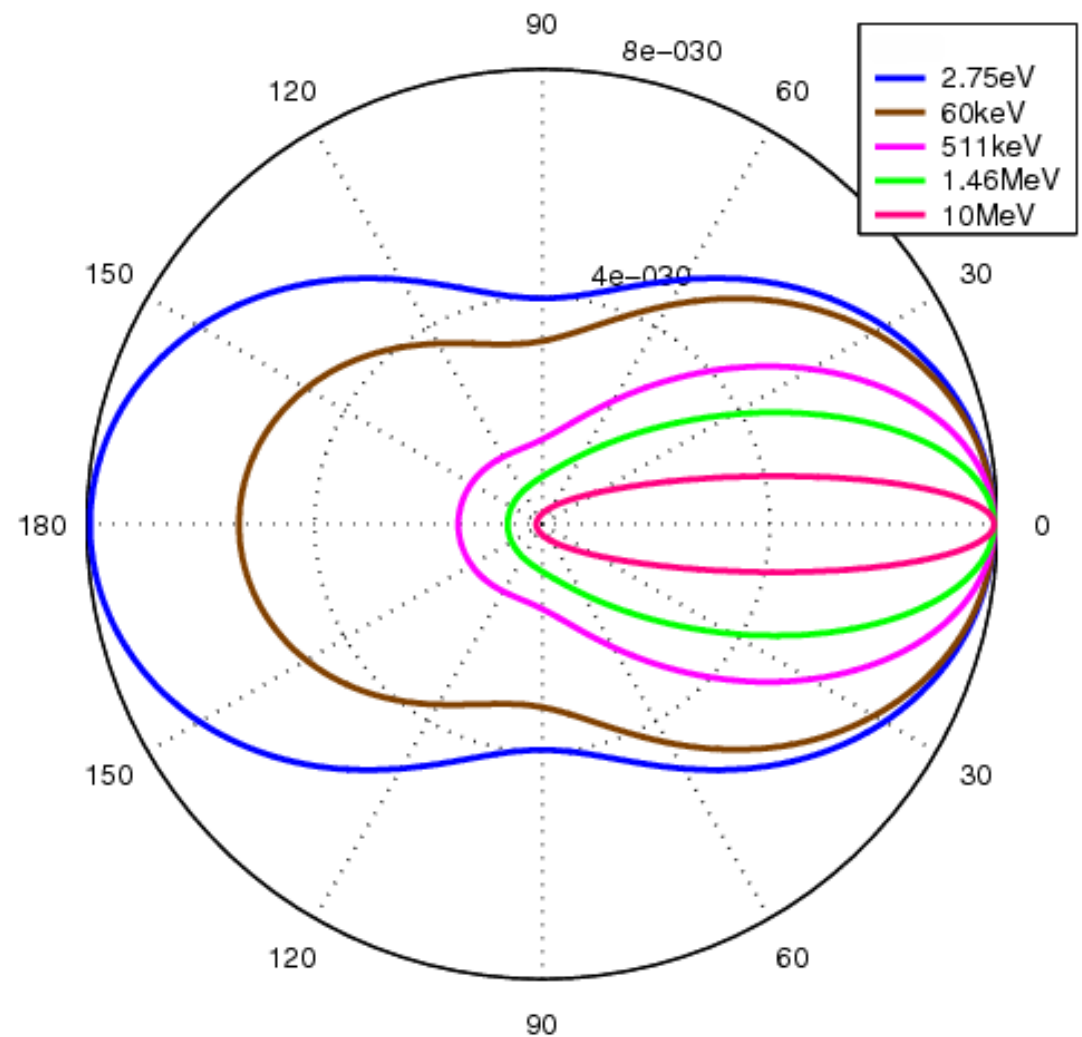
Phase function for scattering-induced polarization

Electron scattering
(Thomson, Compton,
Rayleigh scattering)

Polarization phase function:

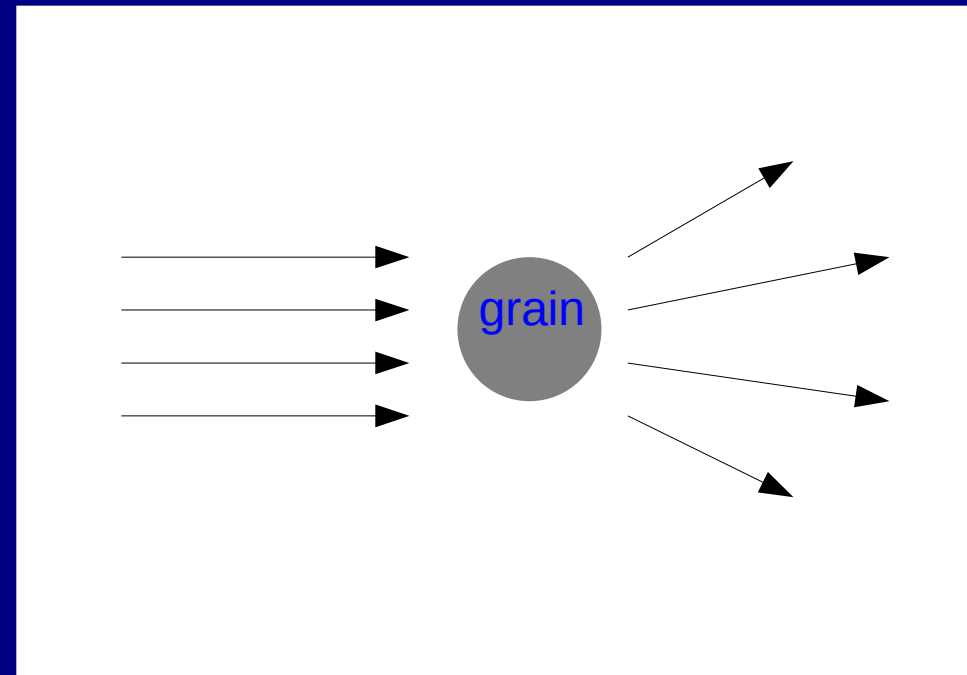


Differential cross section



Polarization by dust scattering

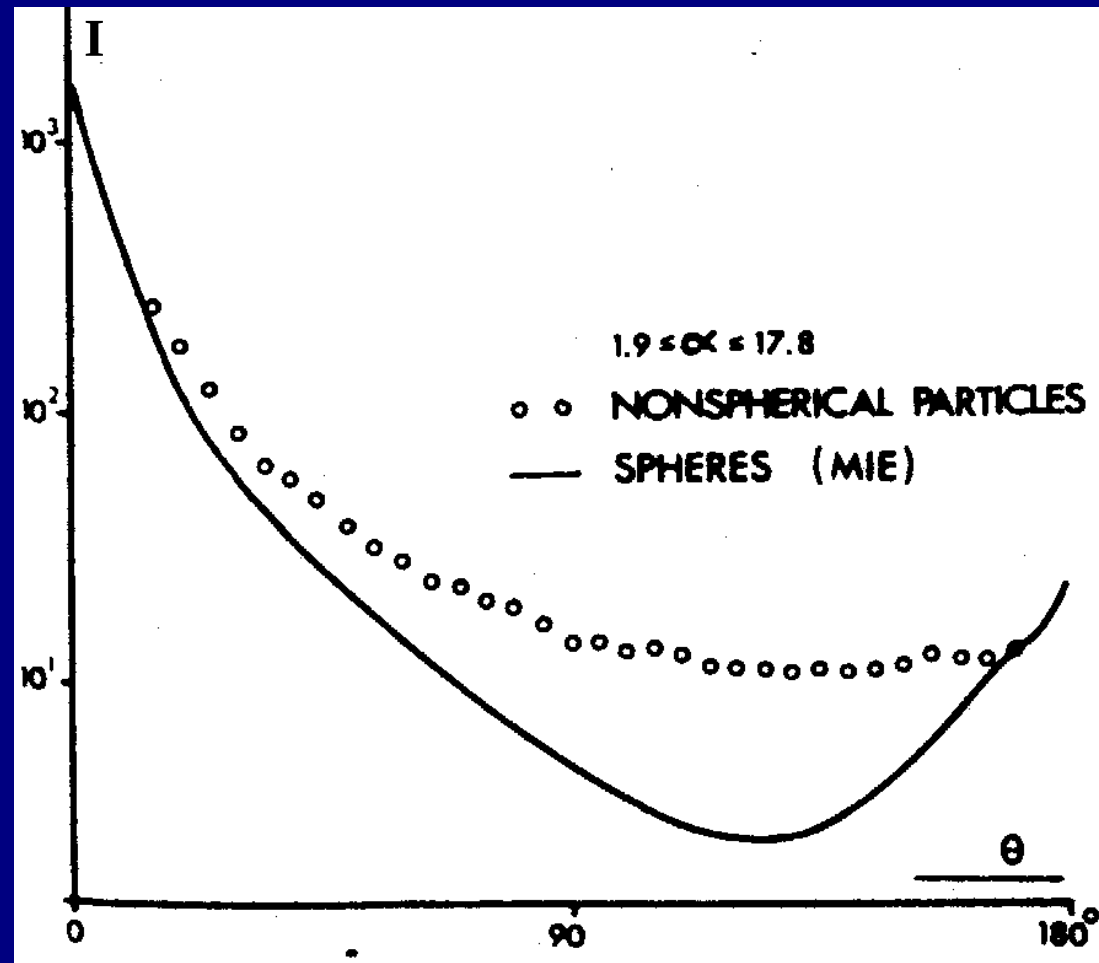
- A plane wave passes by a spherical dust grain.
- The electromagnetic field around the grain is computed by solving the Maxwell equations (Mie theory)
- A complete solution for the scattered radiation including polarization is obtained.



Polarization by dust scattering

It turns out that...

- ...the intensity phase function **favors forward-scattering over back-scattering.**
- ...as for electron scattering, the resulting **polarization is stronger for perpendicular scattering angles.**



Reflection-induced polarization



Reflection-induced polarization



Radio-quiet objects, hidden type-1 AGN

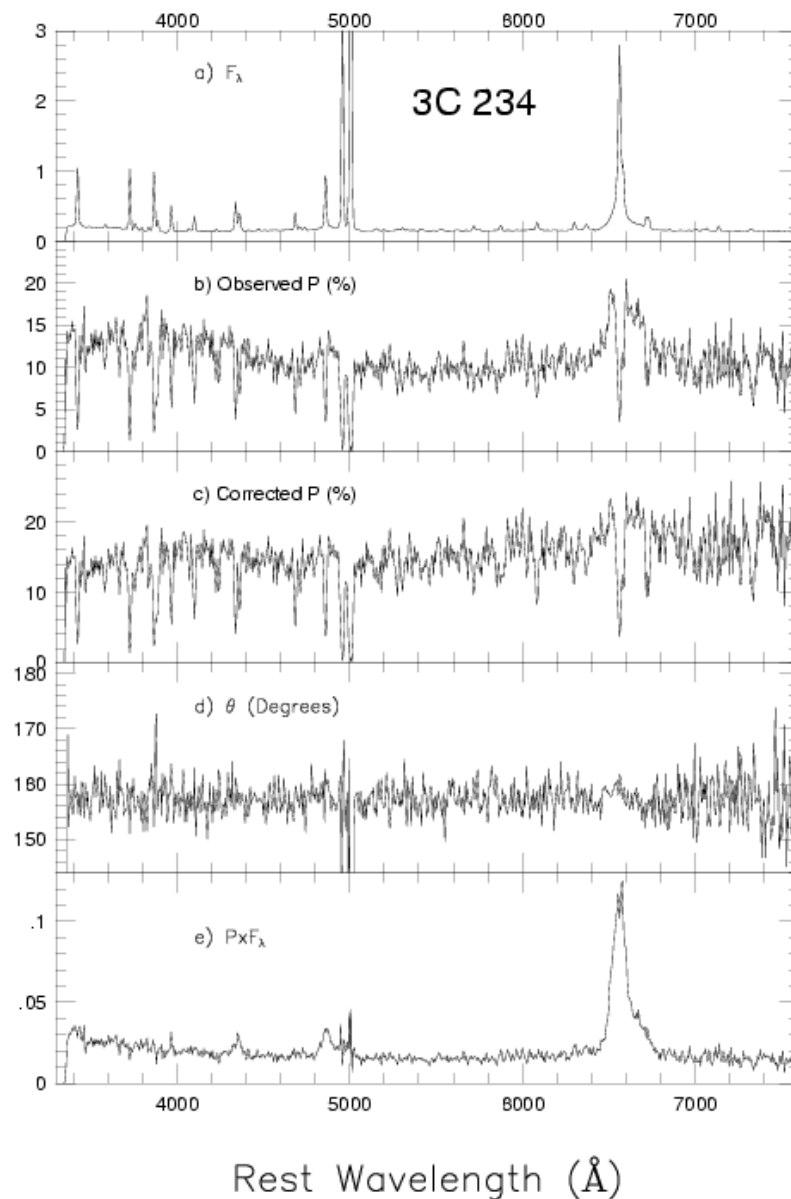
In the following of [Antonucci & Miller \(1985\)](#), more and more hidden type-1 nuclei were found in Seyfert 2 galaxies.

The polarization dichotomy of AGN was established:

Type-2 → P.A. \perp jet axis

Type-1 → P.A. \parallel jet axis,
except for dominant polar scattering

See [Antonucci \(1993\)](#) and [Smith et al. \(2002\)](#) and references therein for summary

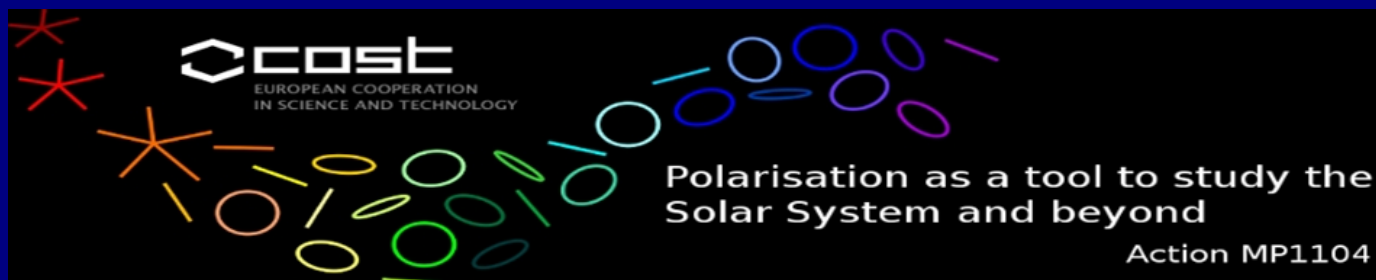


Tran et al. (1995)

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Modeling (scattering-induced) polarization with STOKES

- Monte-Carlo radiative transfer in 3D
- Various geometries for the emission / scattering regions
- polarization due to (multi-)electron scattering and dust (Mie-)scattering
- Resonant line scattering included
- Photo- and K-shell ionization
- variability and evolution of the system

Public access to version 1.0 (and soon 1.3)

<http://www.stokes-program.info/>

home - Mozilla Firefox

Fichier Édition Affichage Historique Marque-pages Outils Aide

http://www.stokes-program.info/

Les plus visités Getting Started Latest BBC Headli... file:///usr/share/appli...

STOKES
Modeling Radiative Transfer and Polarization

Observatoire
Astronomique
de Strasbourg

René W. Goosmann

home
manual
program
examples
dust models
source codes
scientific results
contact

The *STOKES* computer program is a Monte Carlo radiative transfer code for modeling multi-wavelength polarization. It was designed to model astrophysical objects of various geometries and considers polarization induced by electron and dust scattering. If you are interested in polarization and radiative transfer you might want to follow the links on this page to find out more about *STOKES*. The code is freely available for use. We just ask if you publish results based on *STOKES* computations that you refer to the [Goosmann & Gaskell \(2007\)](#) paper describing the code.

--> Check out the [manual](#) to get an introduction to the capabilities of the program and learn how to run it.

--> You may download compiled versions of the [program](#) for Linux and Windows.

--> [Examples](#) for the input files and pre-computed [dust models](#) are helpful to get used to running the code.

--> If you would like to adjust the program to your personal needs and compile it yourself you may obtain the C++ [source codes](#).

--> Find out about [scientific results](#) obtained with *STOKES*.

STOKES was written by **René W. Goosmann** who is now at the Observatoire Astronomique de Strasbourg, France. If you have questions or comments about the code, please [contact](#) him.

George Gabriel Stokes
(1819-1902)

$p \times F$ [arbitrary units]

λ [Angstrom]

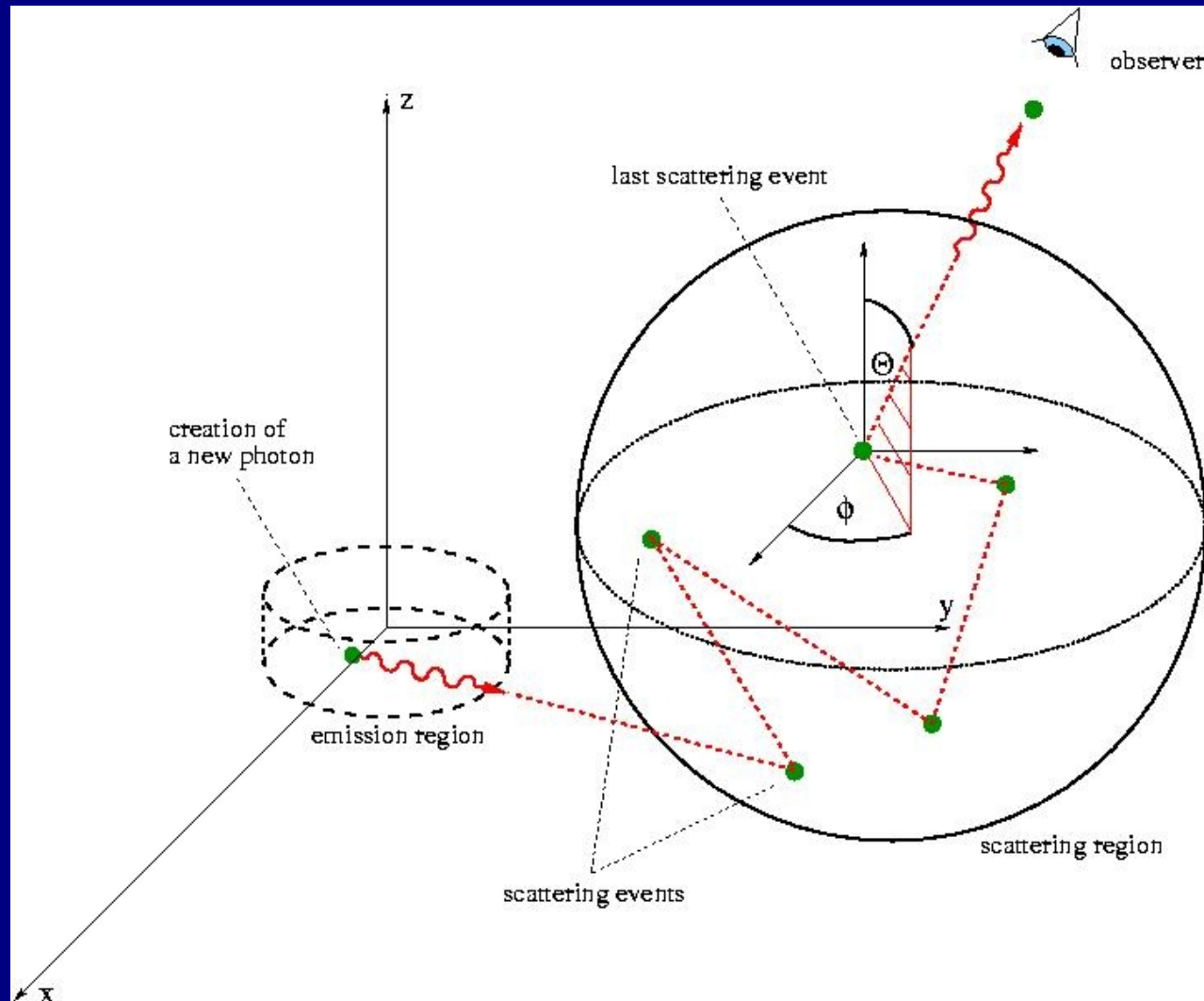
x y z

Terminé

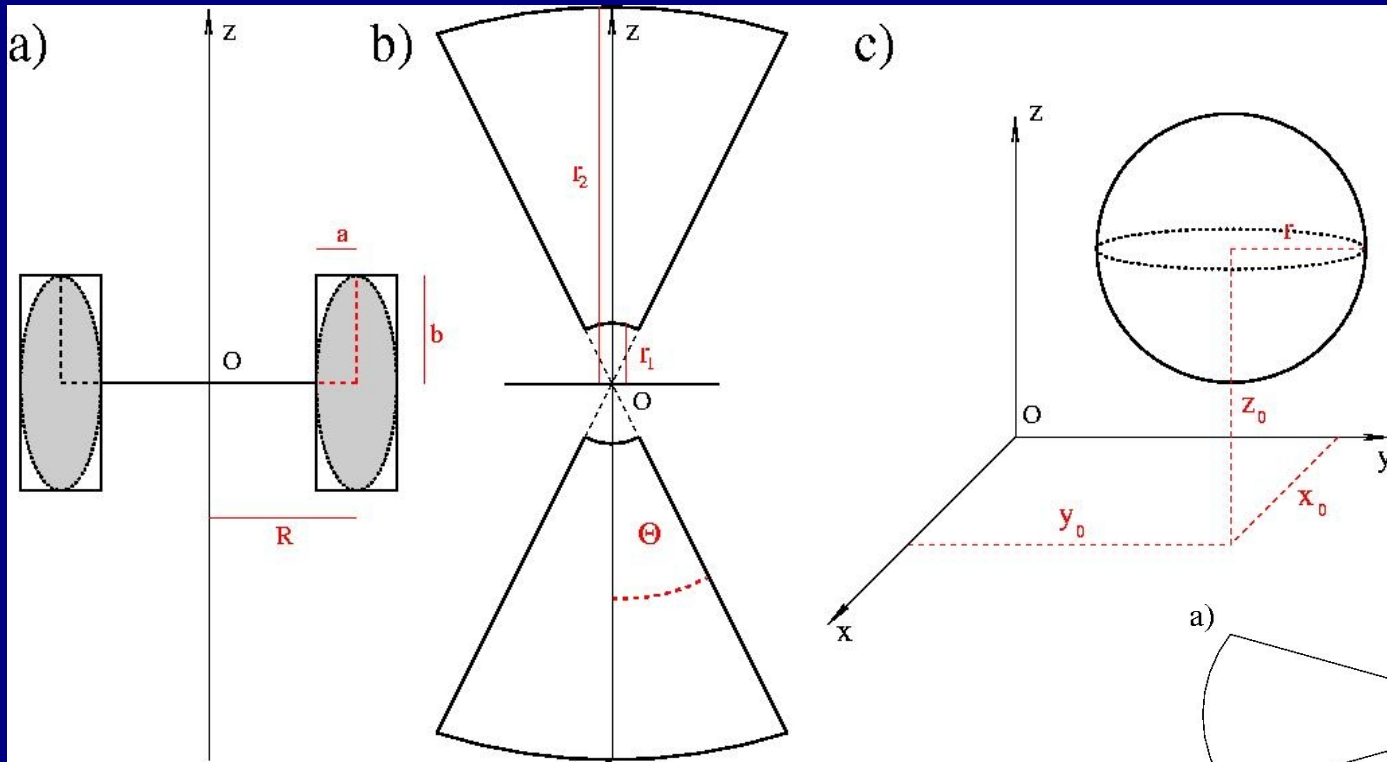
Take a photon for a (random) walk...

Throw the dice for...

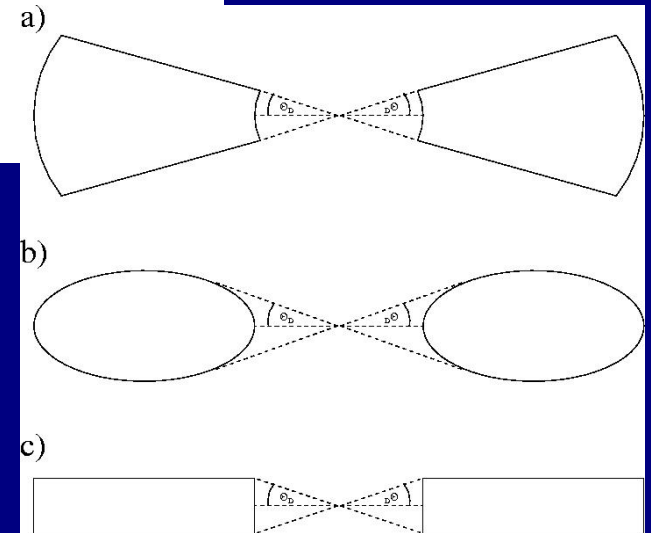
- Photon wavelength
- Emission direction
- Free path length
- Absorption by dust
- Scattering angles



Let it encounter scattering clouds...

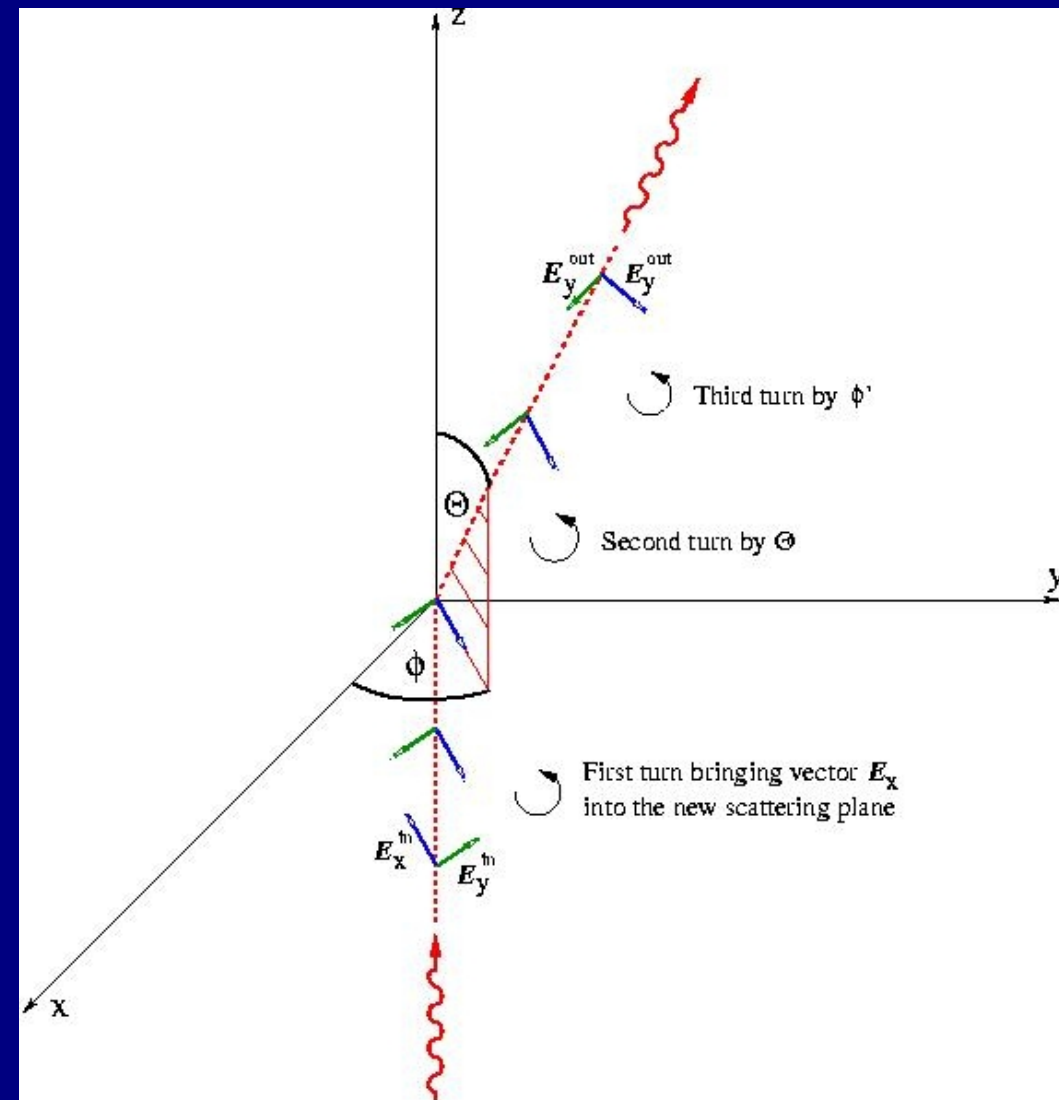


It is possible to define several scattering regions in the model space...



Have the photon scatter accurately...

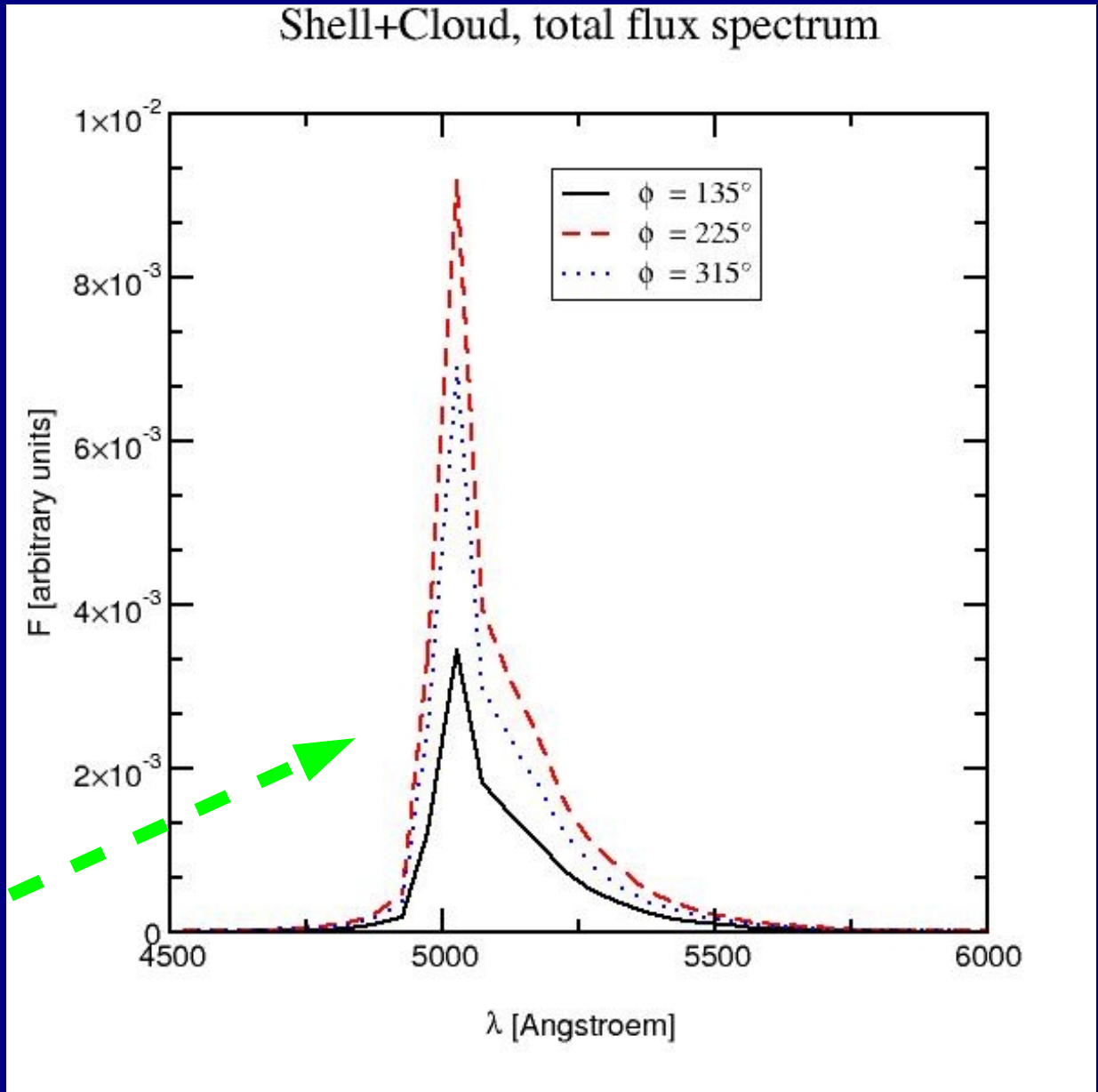
- scattering events are 3D, i.e. around two angles
- the sampling of the scattering angle **depends on the initial polarization state**
- multiple scattering is automatically included
- for dust grains, absorption is considered according to the albedo
- The **correct polarization matrices for resonant line scattering** are taken into account



Count the photons that manage to escape...

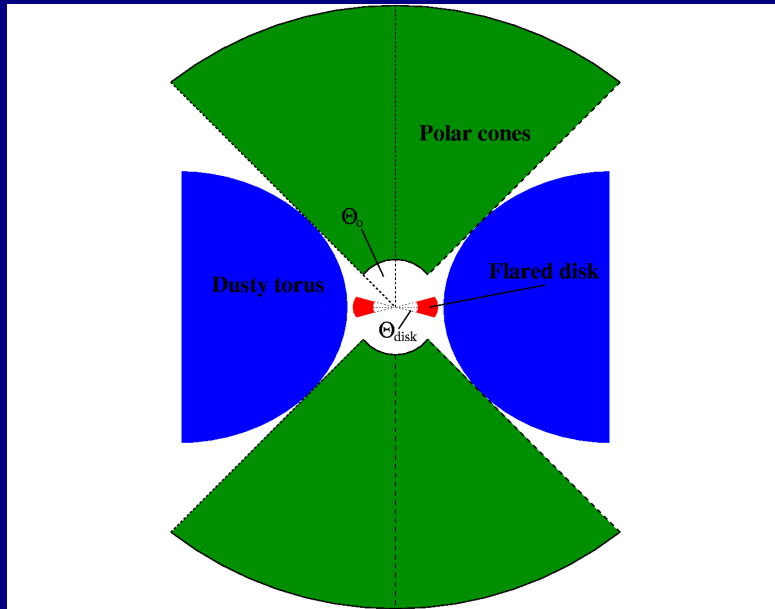
- Stokes parameters are added up for each viewing direction
- number of scatterings are counted
- flight time of the photons is registered

Total flux spectrum of an [OIII] $\lambda 5007$ line for a centrally illuminated, expanding shell with a few dust clouds further out

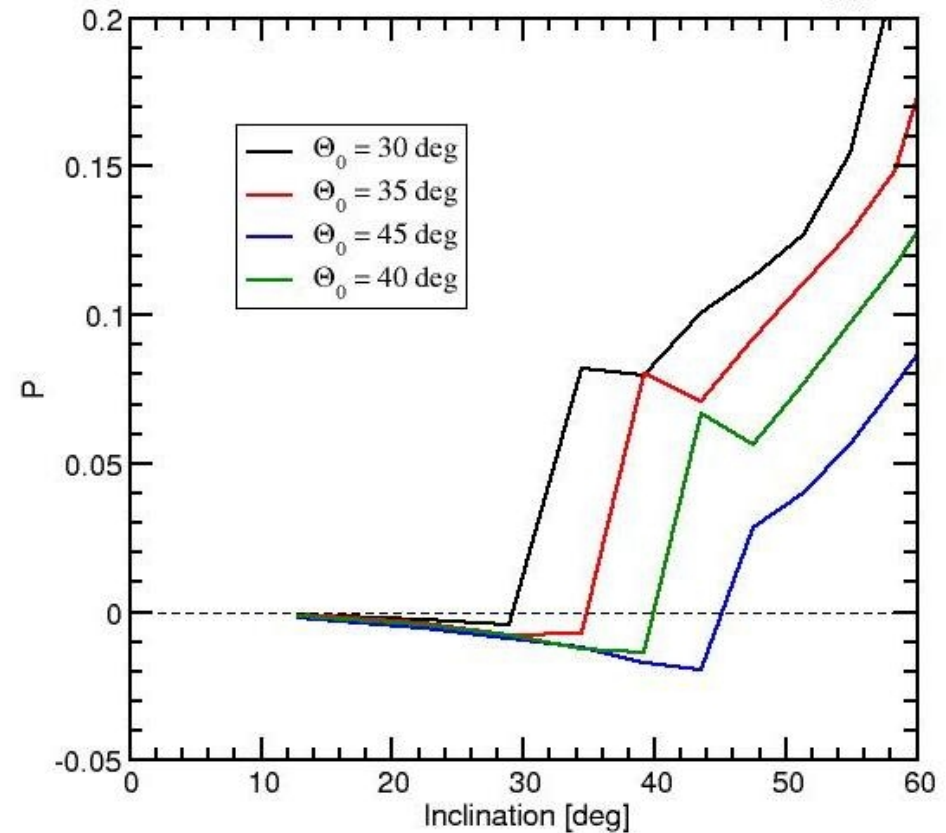


Toward the unified model...

Modeling a composition of **flared electron disk**, **dusty torus**, and **polar electron cones**.



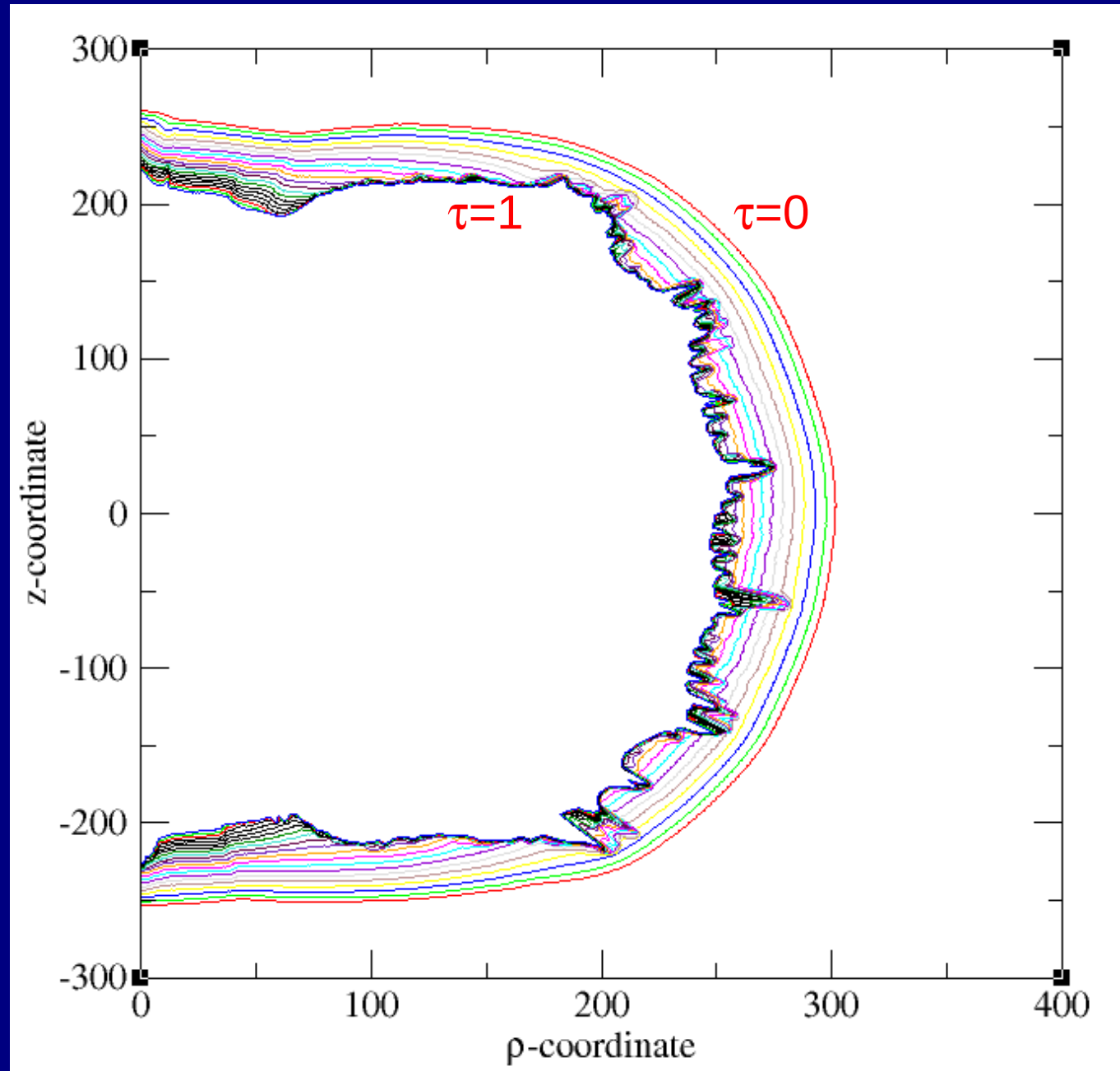
Flared-Torus-Cone, config. with max. type1 polarization
for different half-opening angles of the torus and the cones, $\tau_{\text{cone}} = 0.01$



The scattering regions can be complicated...

Hydrodynamic modeling of supernova explosions

The explosions are anisotropic (due to the rotation of the progenitor)



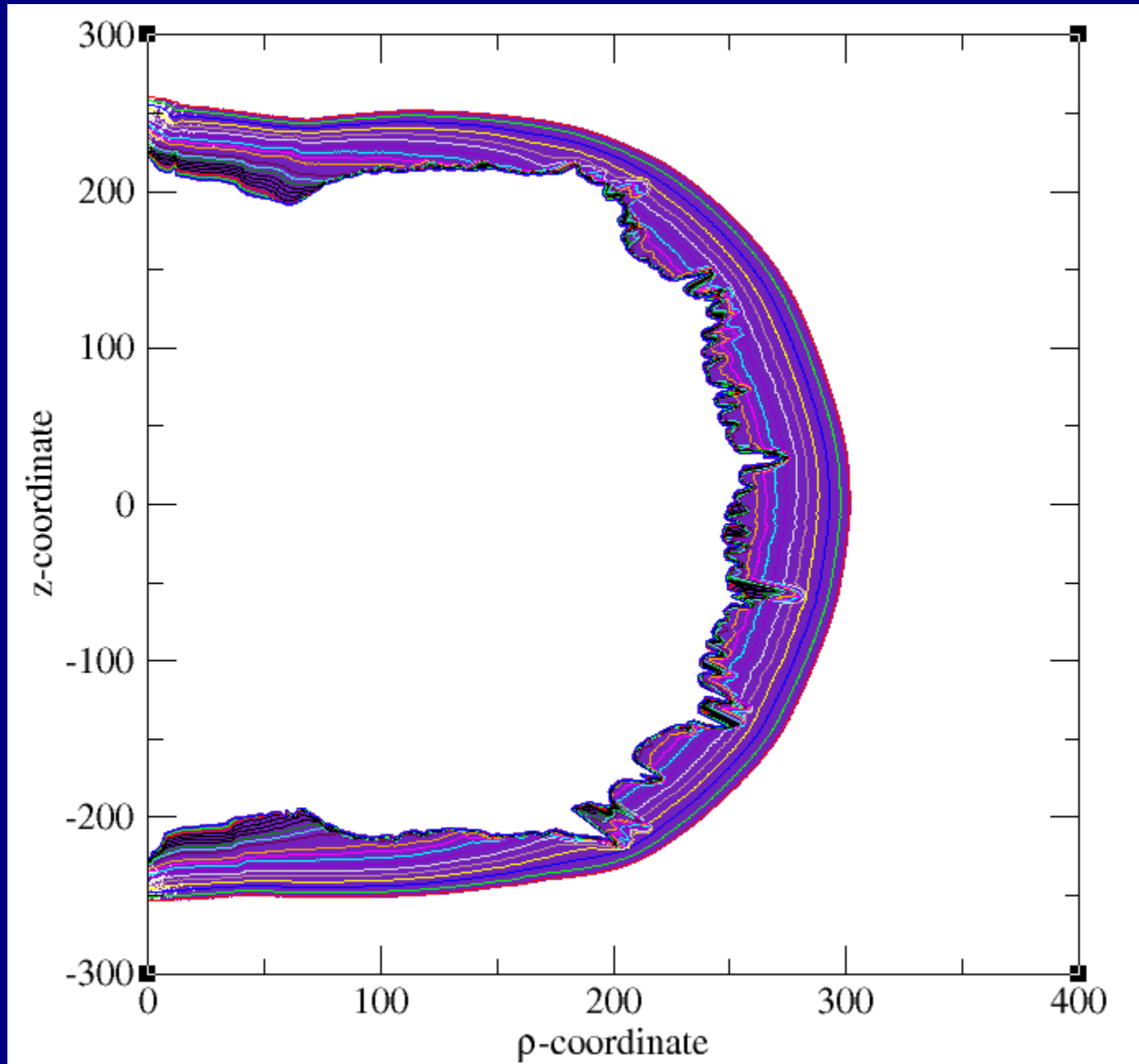
Sean Couch, Craig Wheeler, et al. (University of Texas)

Hydrodynamic simulations of Supernovae

Hydrodynamic modeling of supernova explosions

The explosions are anisotropic (due to the rotation of the progenitor)

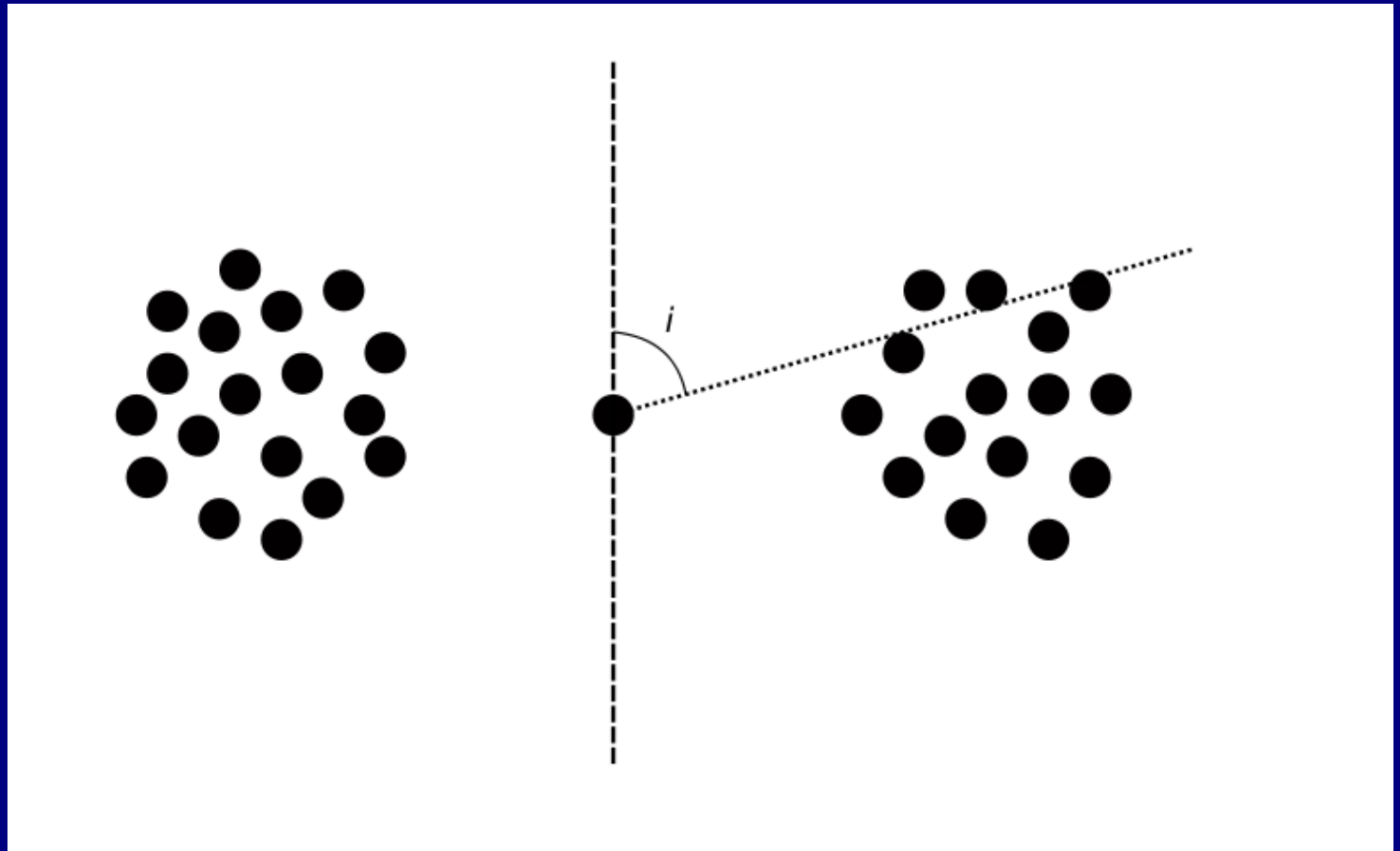
Use *STOKES* to simulate the scattering between optical depths of $\tau=1$ and the photon escape



Sean Couch, Craig Wheeler, et al. (University of Texas)

Towards more sophisticated torus models for AGN

- Polarization models of clumpy tori (e. g. Nenkova et al. 2008, Heymann & Siebenmorgen 2012, Höning & Kishimoto 2010)
- To be explored in the light of the work by M. Stalevski, L. Popovic and collaborators...!



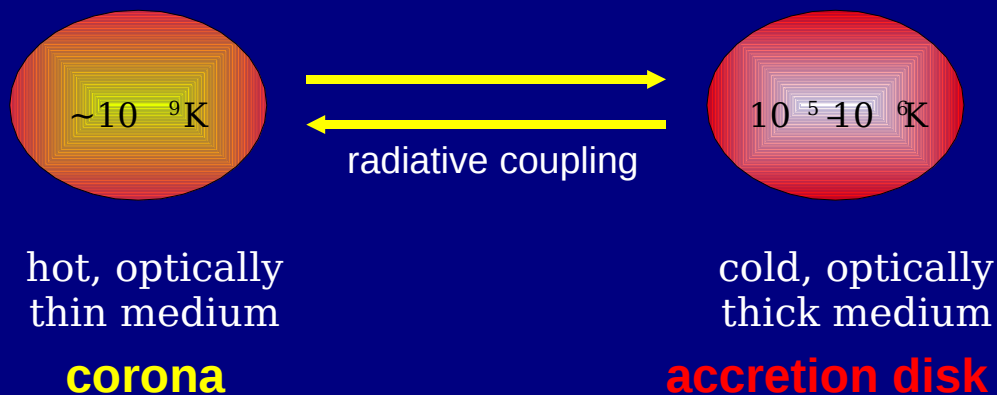
Thesis work by
Frédéric Marin

The X-ray spectrum of AGN

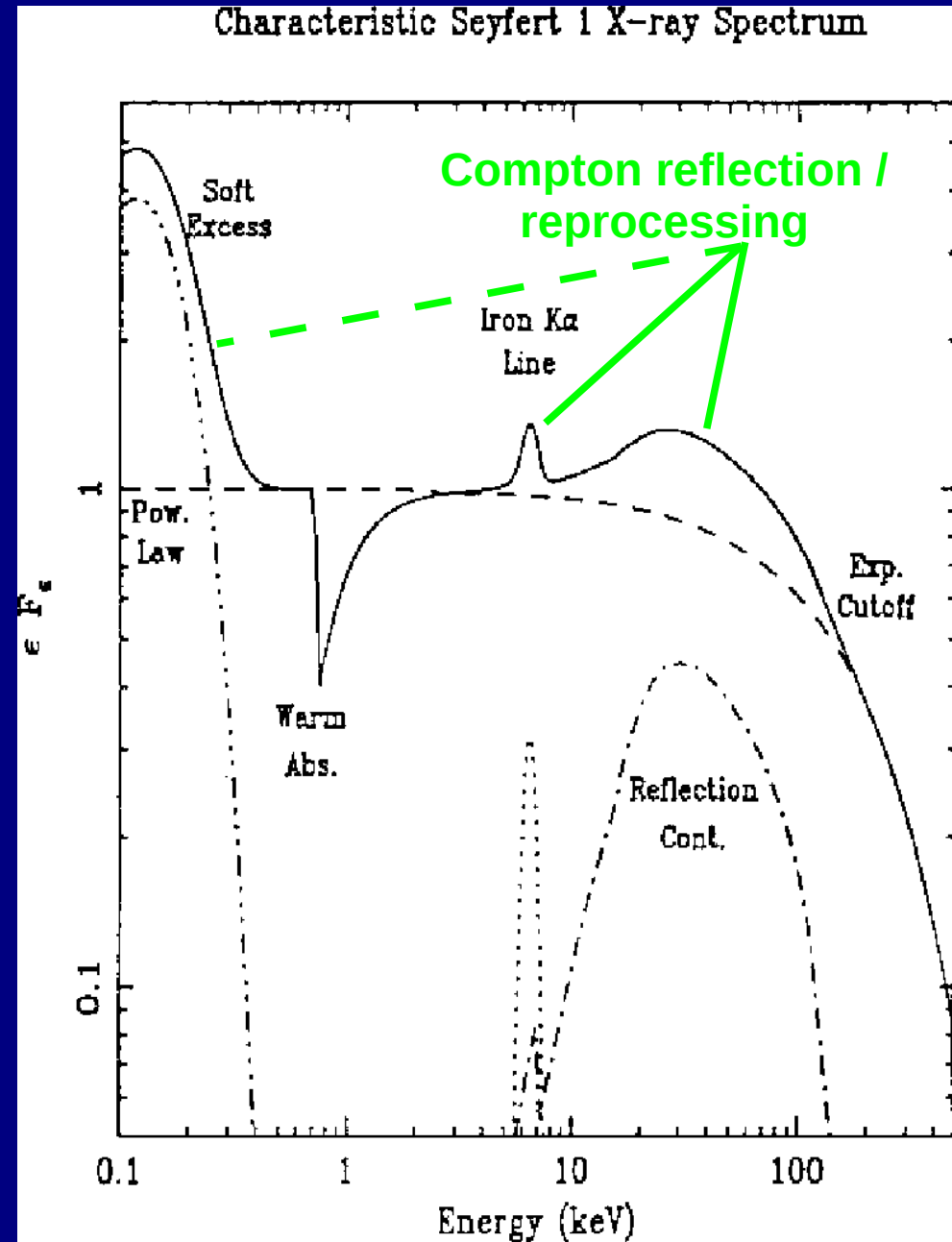
- Primary power-law component
- High-energy cut-off
- Iron $K\alpha$ -line complex
- Compton hump
- Soft X-ray absorption
- Soft-Excess (Crummy et al. 2006)

but: Gierlinski & Done 2004,
Chevallier et al. 2005

Two media are required...

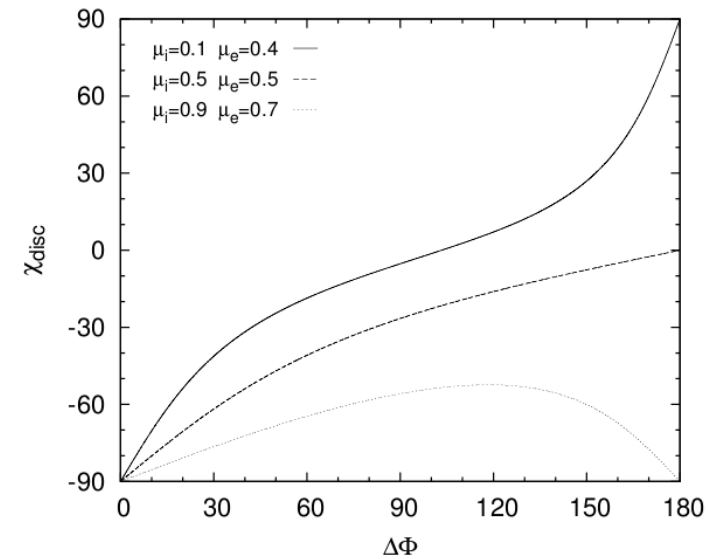
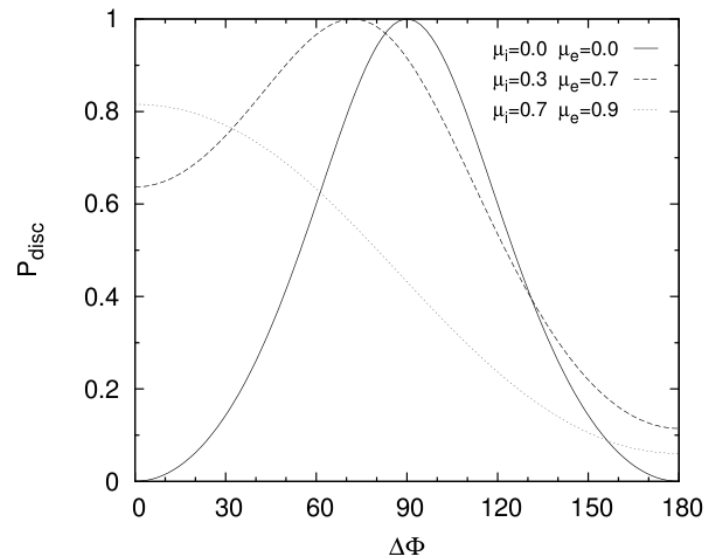
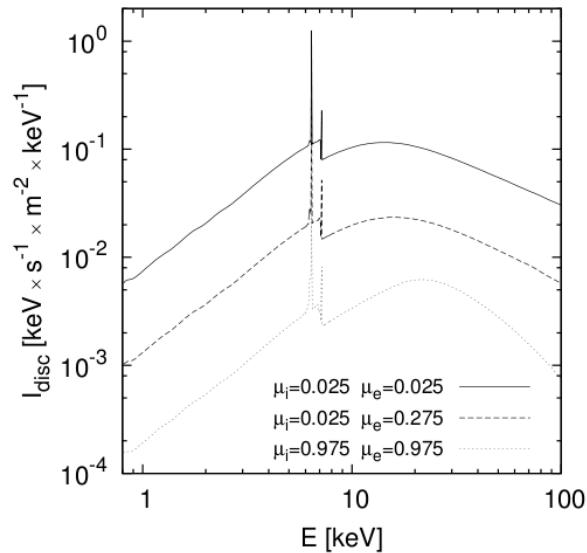


Fabian (2000)

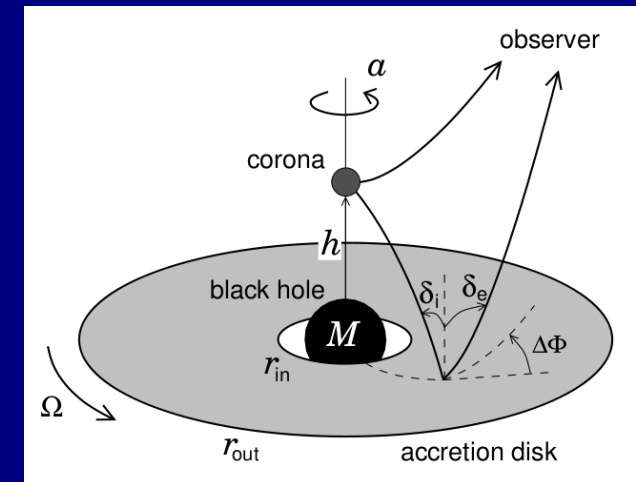


The lamp post geometry

Reprocessed emission in the local disk frame (Dovčiak et al 2011)



- The shape and polarization of the locally reprocessed radiation depends strongly on the position on the disk.
- Polarization degree and angle cover the whole range of possible values.



Including relativistic effects

- Applying relativistic ray-tracing methods in Kerr metric

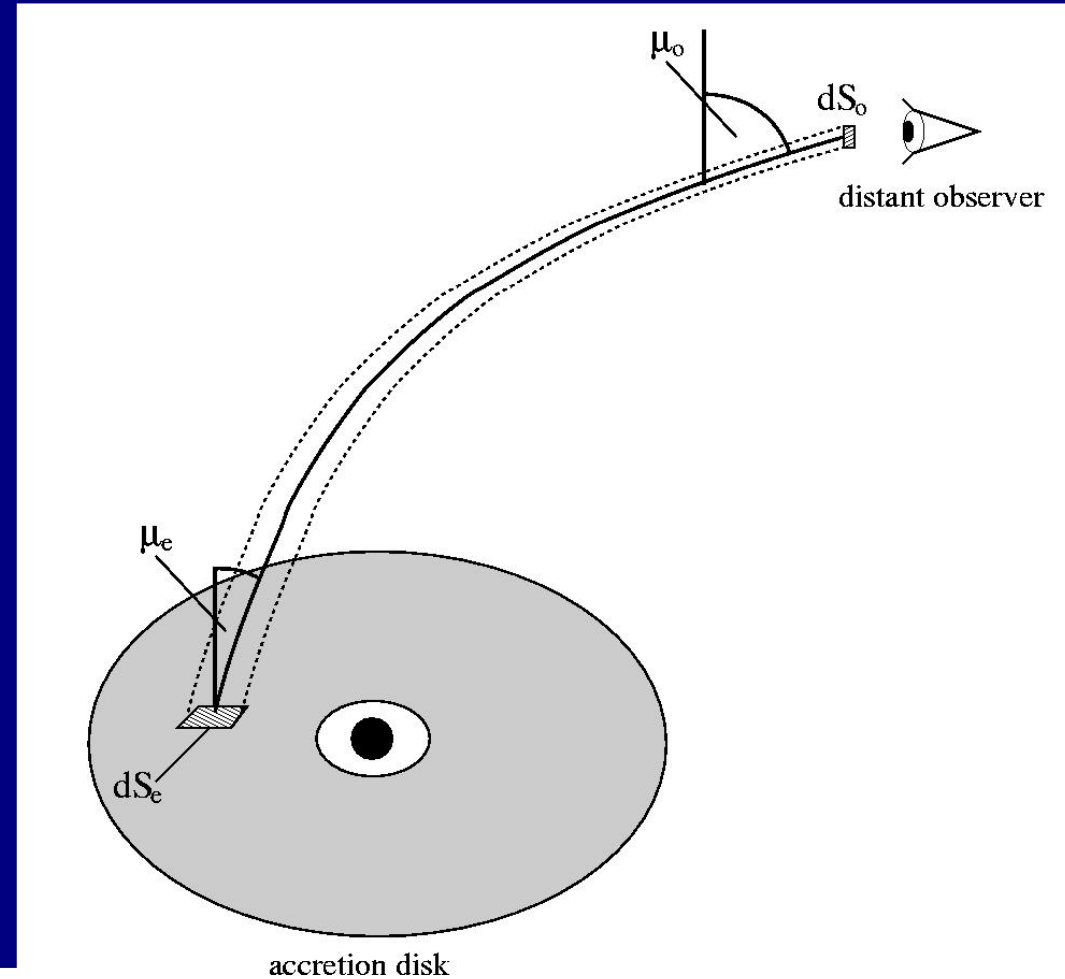
- Important to know the local polarization

see e.g.

Connors, Piran, Stark (1980)

Dovčiak et al. 2004

Schnittman 2009



I, Q, U, V

$$\Delta N_o^{\Omega_o}(E, \Delta E, t) = \int_{r_i}^{r_o} dr \int_{\phi}^{\phi + \Delta \phi} d\phi \int_{E/g}^{(E + \Delta E)/g} dE_l N_l(E_l, r, \phi, \mu_e, t - \Delta t) g^2 l \mu_e r.$$

observed photon flux



disk integration



energy

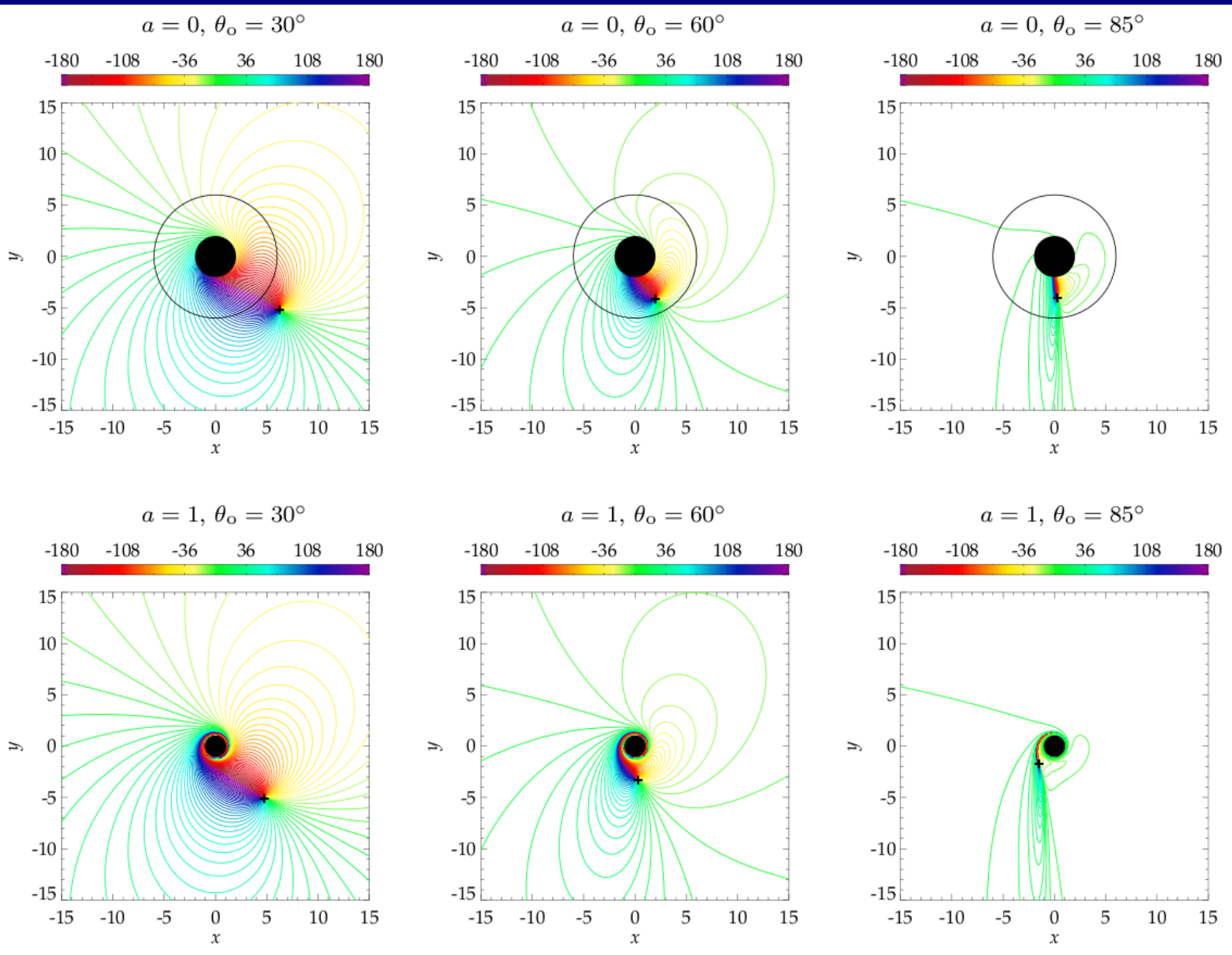


local photon flux



transfer

Integrating the polarization angle



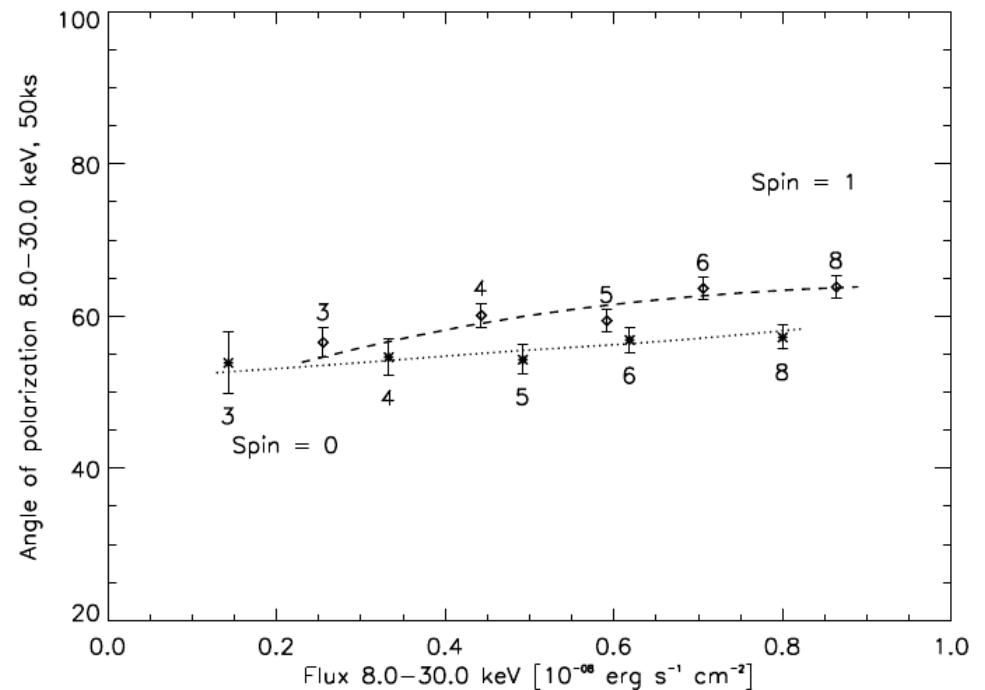
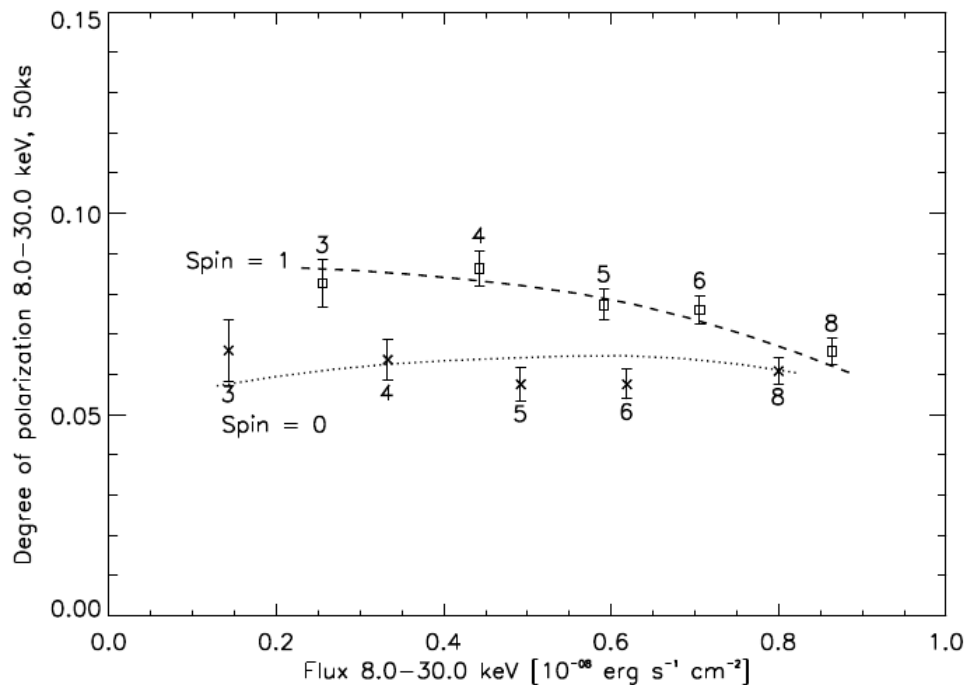
The observed polarization at infinity is obtained by integrating the transferred local polarization.

This gives a vast range in polarization angle...

Moving the lamp post – polarization variability

- Simulation of a series of snapshots for different heights of the primary source (medium size mission including broad-band polarimetry)
- The two extreme spin states are still distinguishable, but less than for the thermal disk

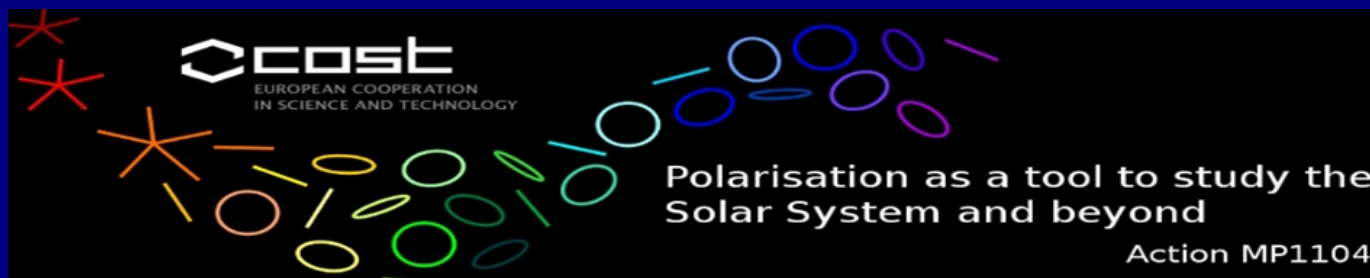
XTE 1650+500 — NHXM simulation of $T = 50$ ksec snapshots, $i = 30^\circ$



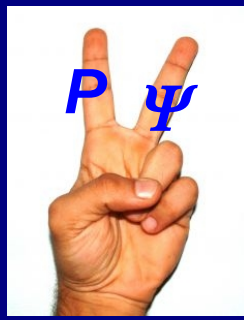
Exploiting and modeling polarization in astrophysics

- The basic concept of (linear) polarization
- Techniques to observe broad-band polarization
- Mechanisms producing polarization
- Modeling scattering-induced polarization with *STOKES*
- **Summary and conclusions**

*Talk by René W. Goosmann at the
9th Serbian Conference on Spectral Line Shapes in Astrophysics
16th May 2013*



SOME CONCLUSIONS...



- Scattering-induced polarization as a function of wavelength reveals details of the interaction between radiation and matter !
- In my opinion, spectropolarimetry still is a neglected source of information about unresolved objects.
- The development of the STOKES code is ongoing and guided by the input of the user – feel free to try it out!

<http://www.stokes-program.info/>

- Some more STOKES results on spectral line are going to come in our talks of tomorrow morning :
 - Line and continuum polarization variability in the optical/UV
 - Absorption, reflection, relativistic effects, and Compton scattering of $K\alpha$ lines in different geometries