

# BH MASS ESTIMATION

## for Quasars

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# Estimating BH Mass

Broad Emission Lines as virial estimators

FWHM Hbeta  $\rightarrow$  virial  $v$

Single epoch FWHM measures? rms? Composites?  
Or sigma?

$$M_{\text{BH}} = \frac{v^2 r_{\text{BLR}}}{G}$$

$$v = f_{\text{FWHM}}$$

$$v = \sqrt{3}/2 \text{FWHM}(\Pi\beta_{\text{BC}})$$

# The f factor?

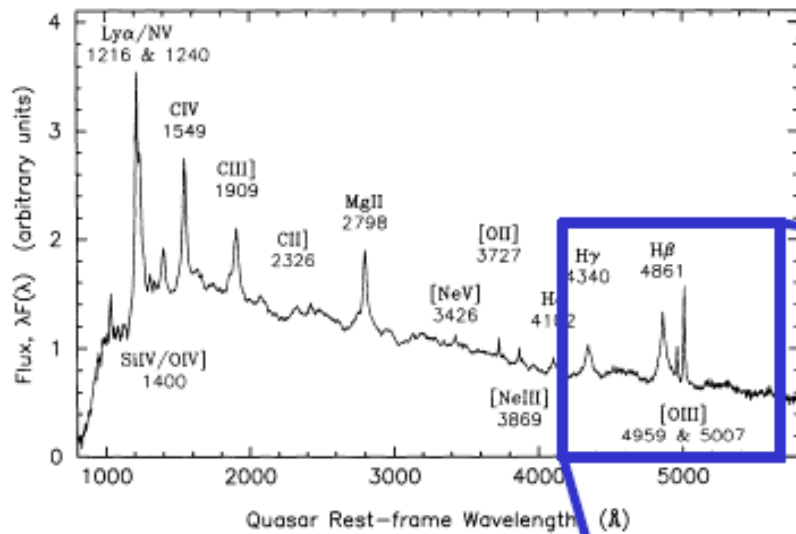
- BLR geometry and kinematics
- Strong inclination dependence
- Of order  $\sim 1-2$  (or 5.5 Onken et al. 2009 using sigma)
- Same for Pop. A and B?

# Not all FWHM Hbeta are the same

Sources above and below FWHM  
Hbeta=4000km/s show different  
structure (Sulentic et al. 2002)

Relationship between FWHM and sigma changes --  
reflecting multiple components

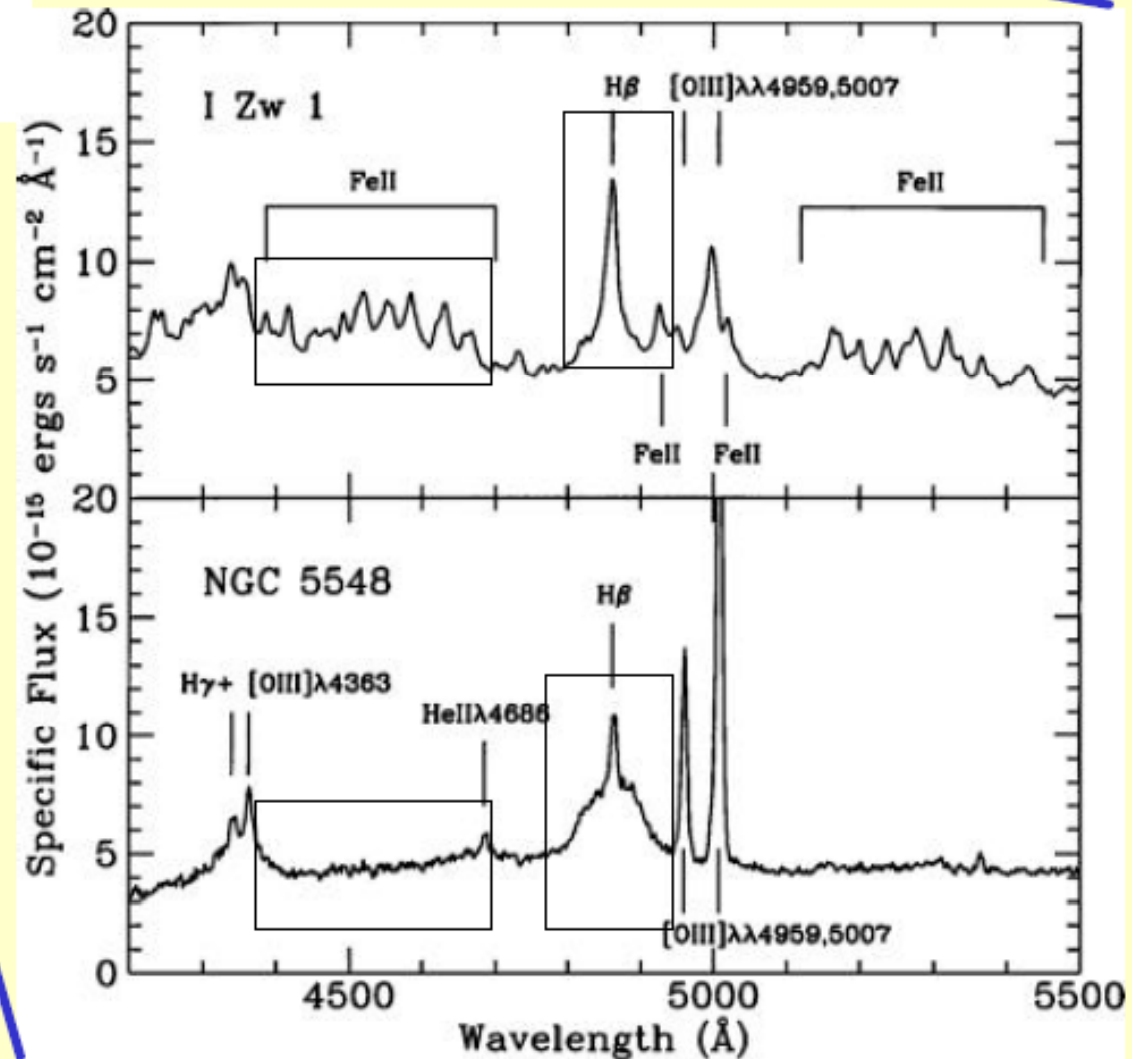
(Collin et al. 2006:  $f(A) \approx 2.12$   $f(B) \approx 0.5-1.0$ )

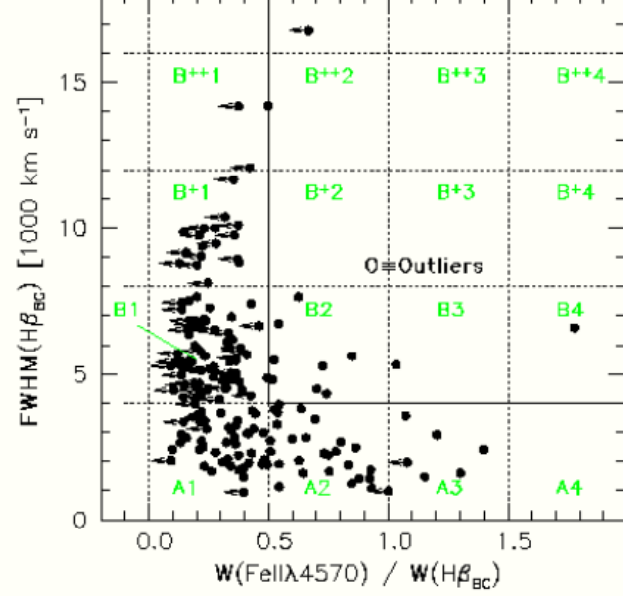


# “Eigenvectors of Quasars”

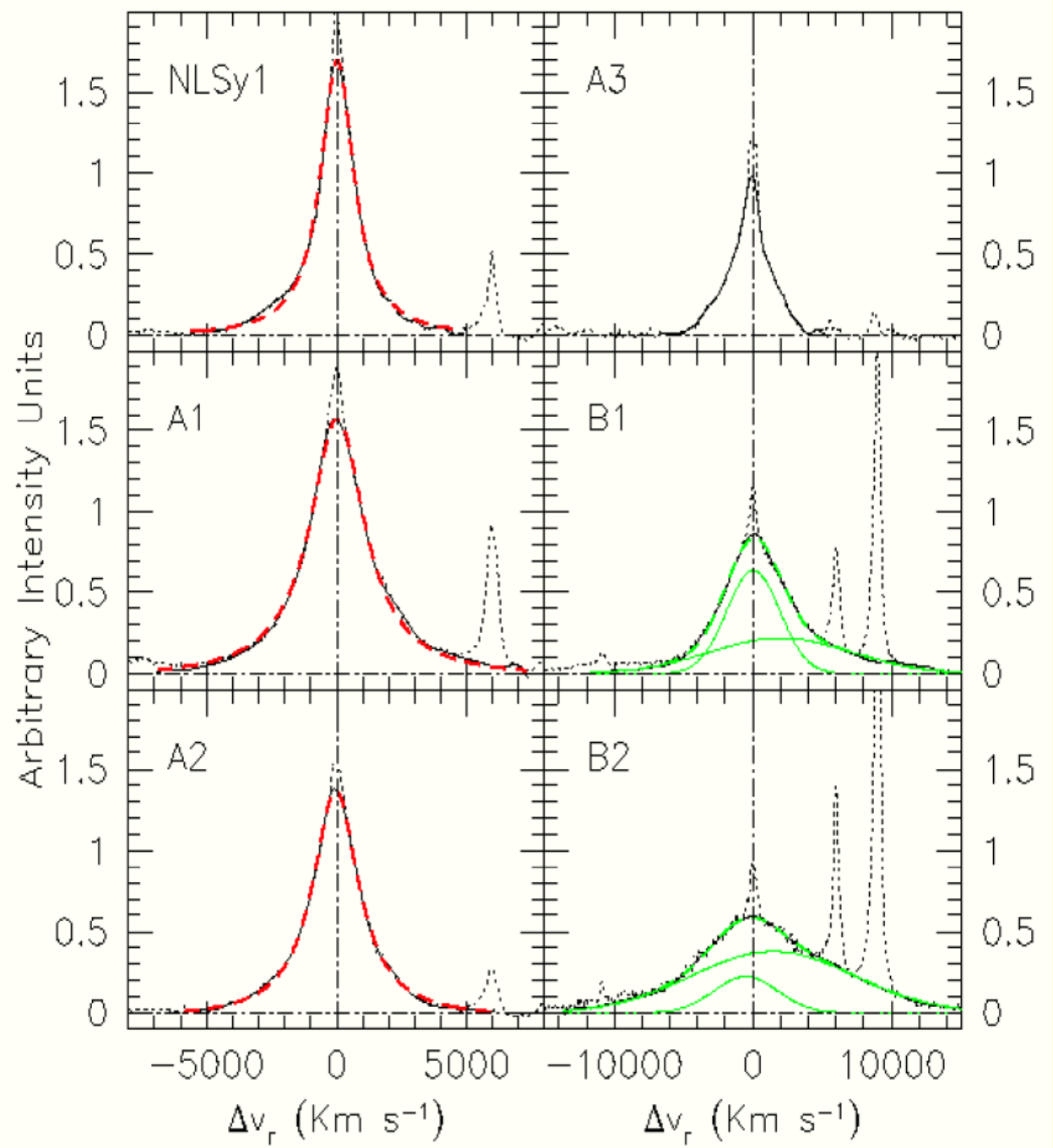
✿ Quasar spectra are not similar!

↳ historic inability to explain line ratios via simple photoionization models





**Sulentic et al. 2002:**  
**LIL/BLR Structural**  
**Difference between**  
**Population A and B**  
 (Sample of about 200  
 Seyfert 1 and  
 low-redshift quasars  
 High S/N and resolution  
 4 Å FWHM)



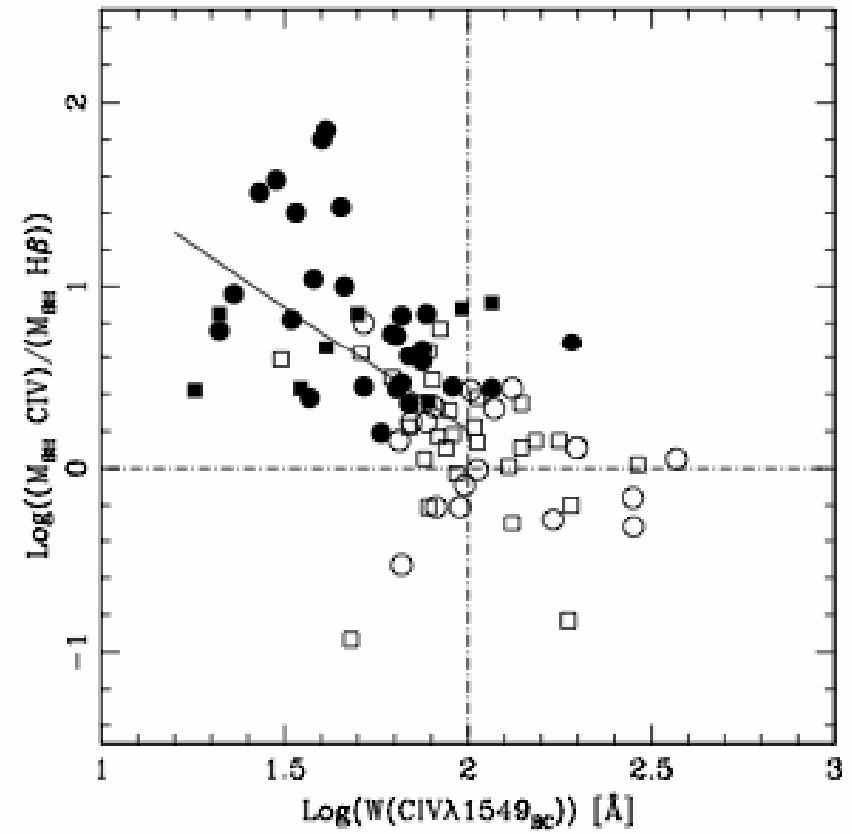
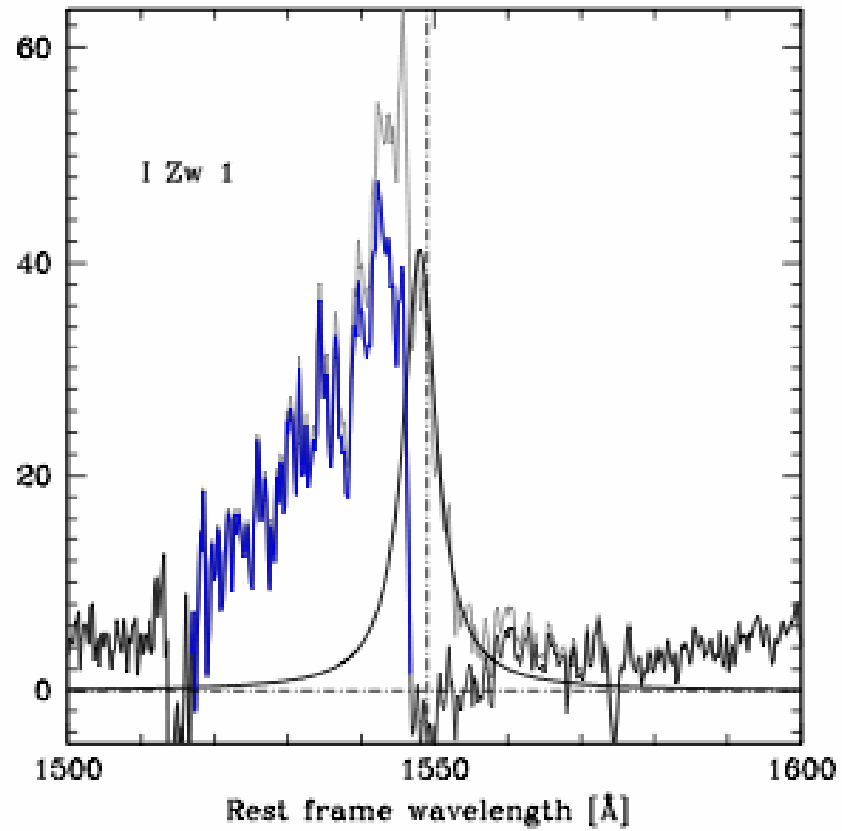
# Extreme $M_{BH}$

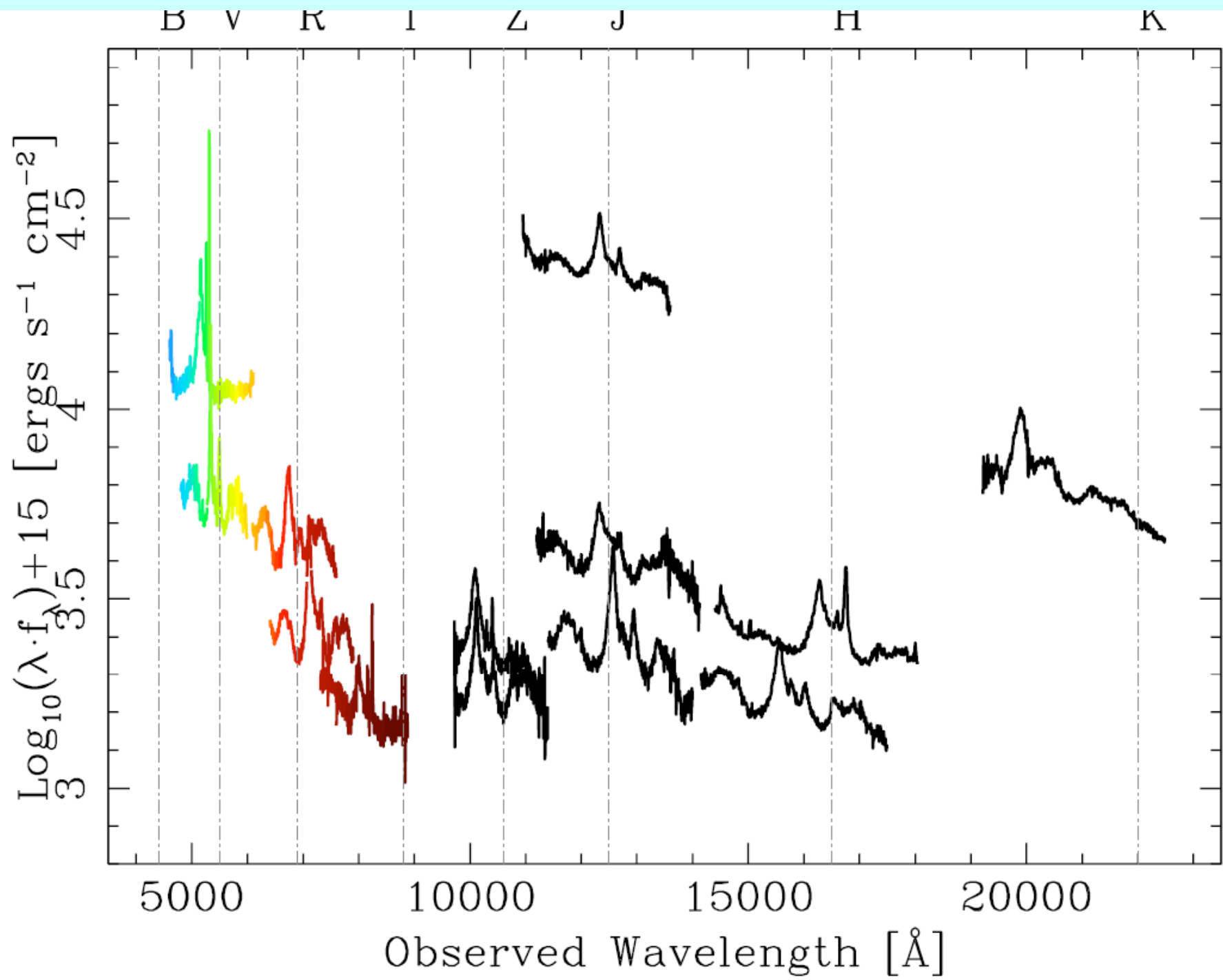
- Pop. A extremely narrow profiles lead to underestimates of  $M_{BH}$
- Most extreme  $M_{BH}$  values came from pop B sources using FWHM Hbeta uncorrected for the extra very broad component.

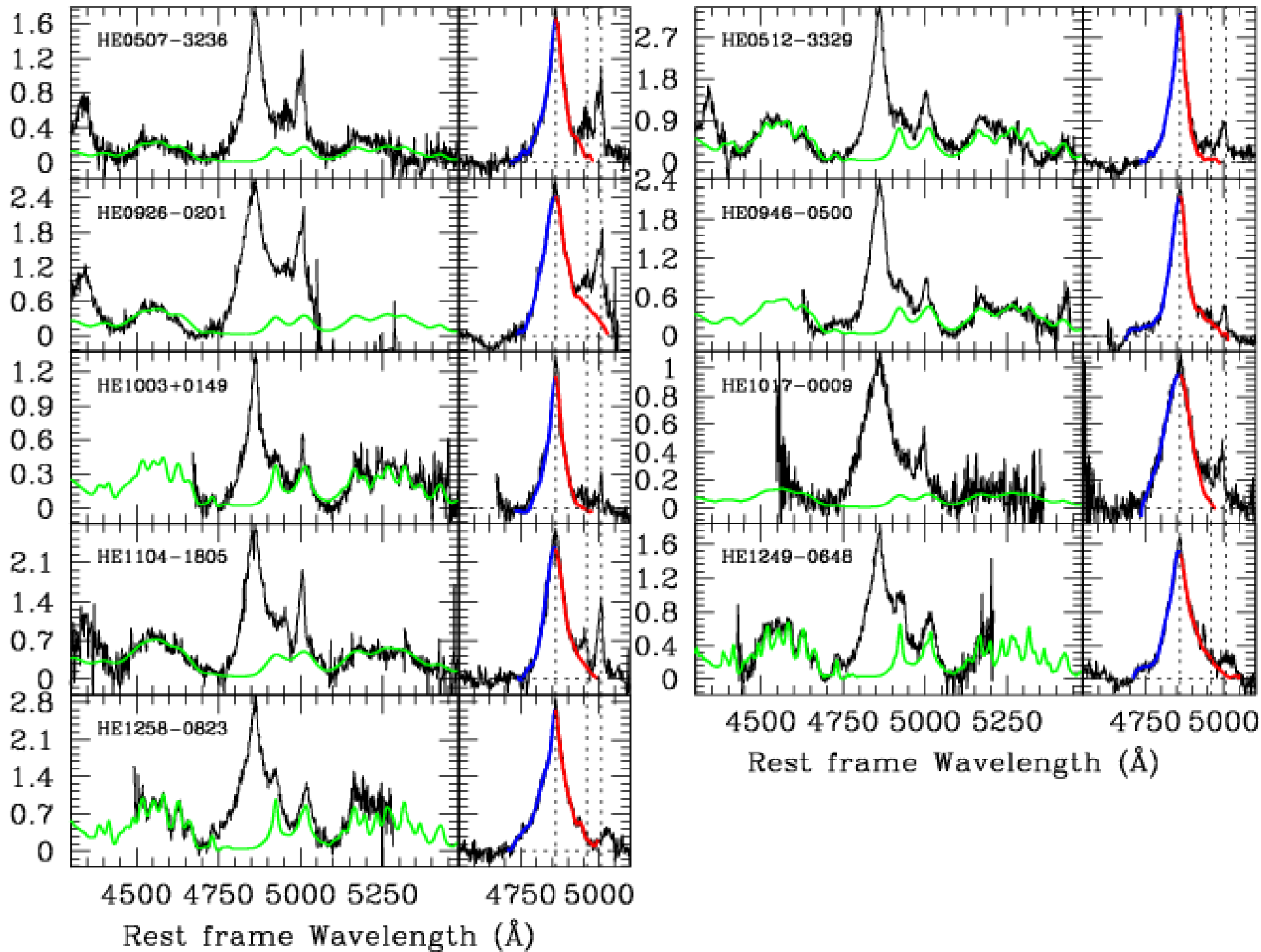


# Above $z \sim 0.7$ ?

- Follow Hbeta into the IR to  $z = 3.8$
- MgII2798 with suitable corrections to  $z > 6$
- CIV1549 dangerous







# So we have “v” – ways to estimate BLR size

- Directly: reverberation mapping

(Peterson & Horne 2006)

- Using  $r_{\text{BLR}}$  vs. Luminosity relation

(Marziani et al. 2009; Trachtenbrot et al. 2011)

- Photoionization Method

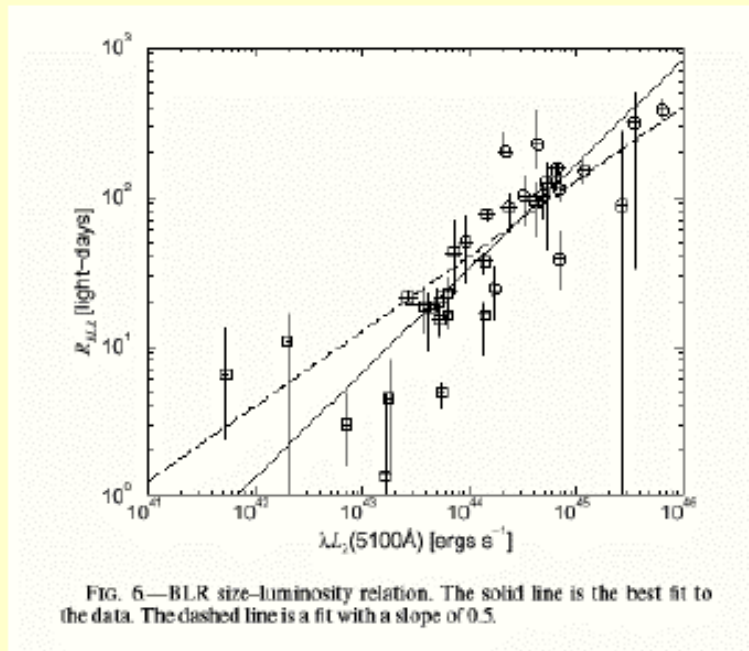
(Dibai 1984; Wandel & Yahil 1985; Padovani & Rafanelli 1989;  
Wandel et al. 1999; Negrete et al. 2011)

## Determination of central compact object mass

$c\tau_{\max}$  provides an emissivity weighted estimate of the BLR linear distance (size) from central continuum source

$$r_{BLR} \approx c\tau_{\max} \approx 33 \left( \frac{\lambda L_{5100}}{10^{44} \text{ erg s}^{-1}} \right)^{0.7} \text{ l.d.}$$

Kaspi et al. 2000



More recent work indicates exponent  $\approx 0.5$

Bentz et al. 2009

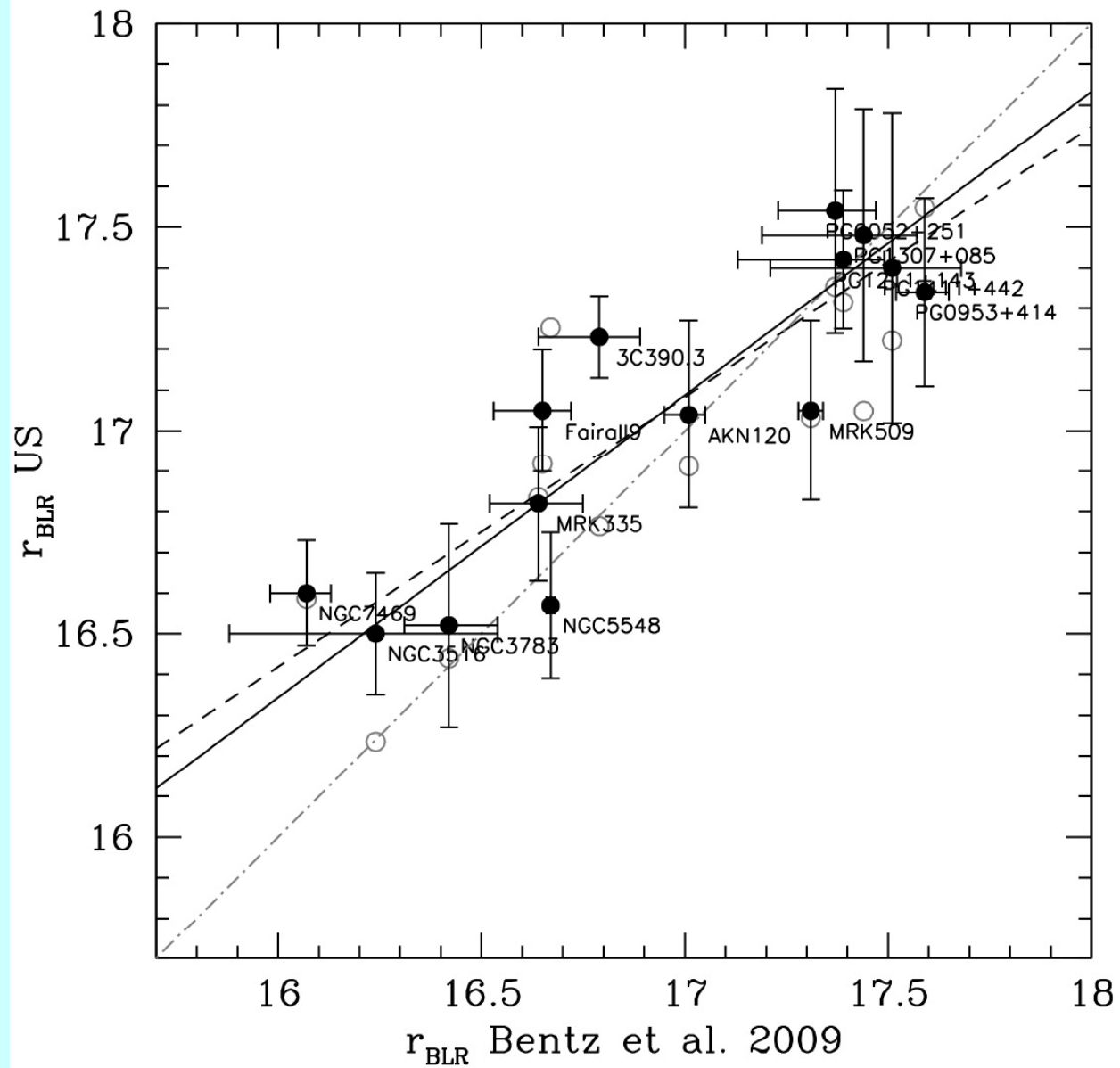
Paves the road to easy mass estimates

Bolometric luminosity estimated from bolometric correction  $\approx 9 - 10$  times the optical luminosity (at 5100 Å)

Elvis et al. 1994; see Nemmen & Brotherton 2010 for a more modern approach

$$U = \frac{\int_{\nu_0}^{+\infty} \frac{L_\nu}{h\nu} d\nu}{4\pi n_H c r^2}$$

$$r_{\text{BLR}} = \underbrace{\frac{1}{(4\pi c)^{\frac{1}{2}}}}_{\text{const.}} \underbrace{(U n_H)^{-\frac{1}{2}}}_{\text{diagnostics}} \left( \underbrace{\int_{\nu_0}^{+\infty} \frac{L_\nu}{h\nu} d\nu}_{\text{\# ionizing photons}} \right)^{\frac{1}{2}}$$





$$M_{\text{BH}} = \frac{3}{4} \frac{r_{\text{BLR}} \text{FWHM}(\text{H}\beta_{\text{BC}})^2}{G}$$

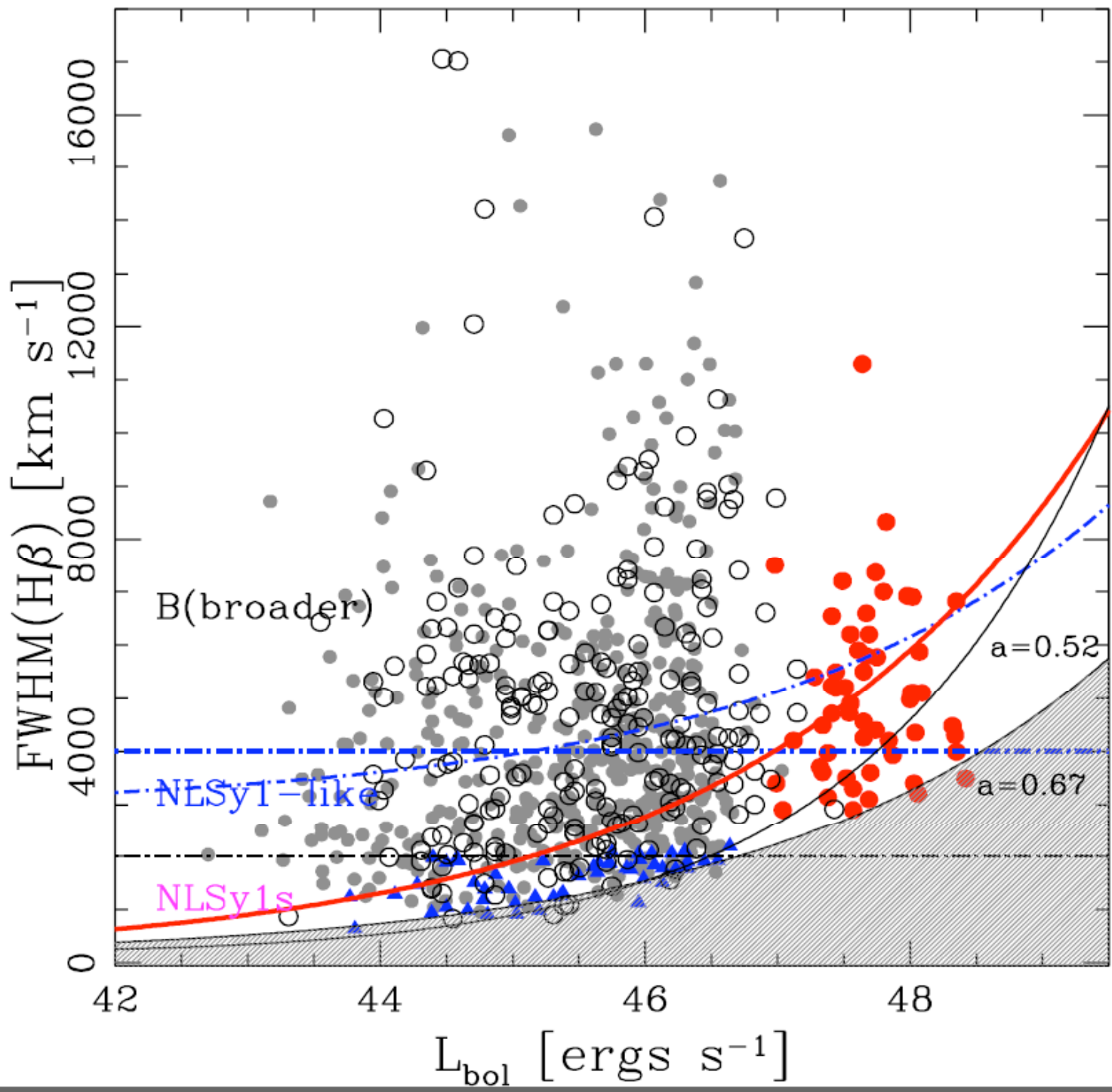
$$r_{\text{BLR}} \propto (L_{5100})^{\alpha}$$

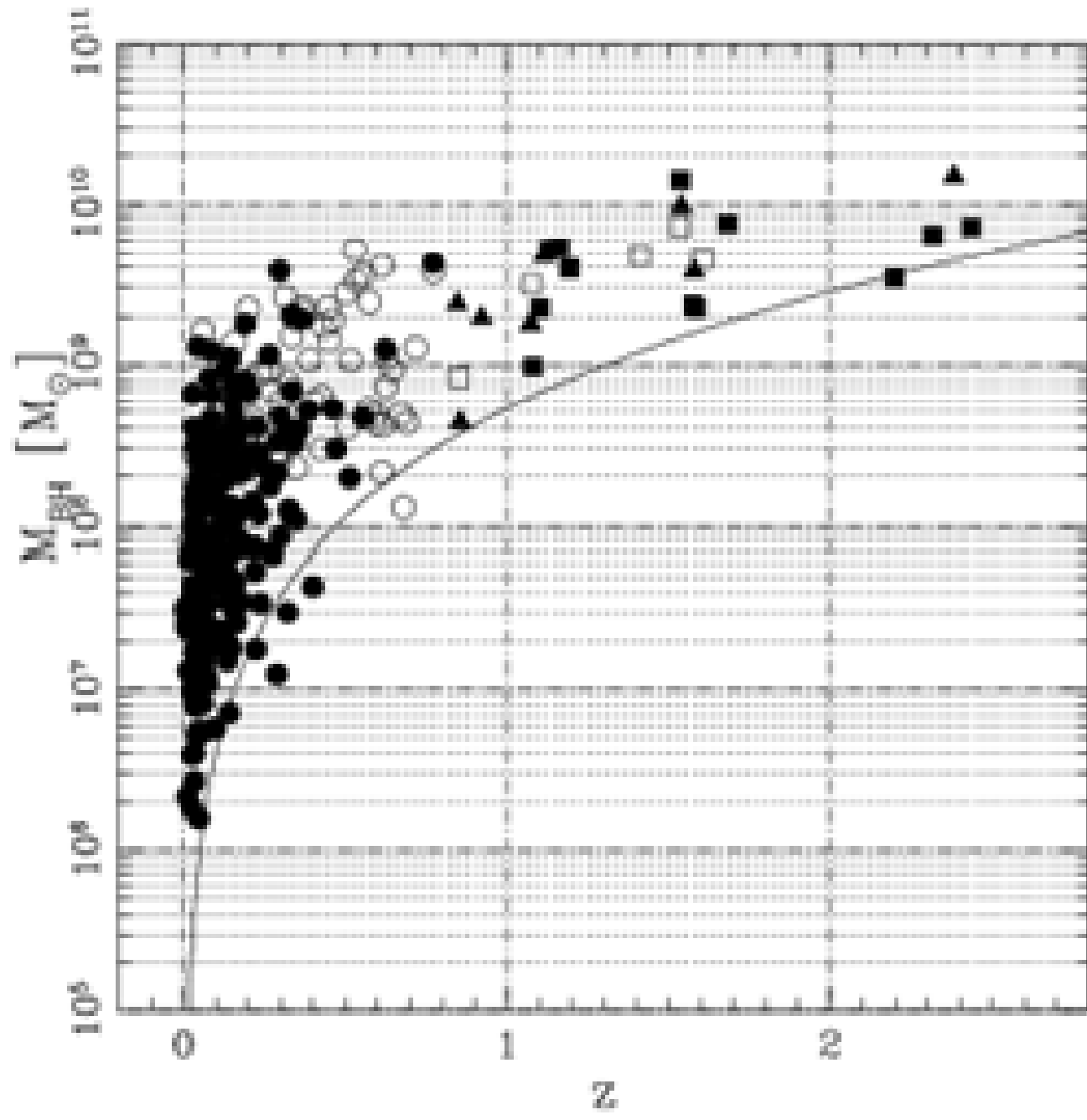
*(Kaspi et al., 2000; 2005)*

$$\tau_{BLR} = 0.85 \cdot 10^{17} \cdot \left[ \frac{\lambda L_{\lambda} (5100\text{\AA})}{10^{44} \text{ erg s}^{-1}} \right]^{0.7} \text{ cm}$$

$$\lambda L_{\lambda} (5100\text{\AA}) = 3.14 \cdot 10^{35-0.4(M_B)}$$

$$\lambda L_{\lambda} (5100\text{\AA}) = 4\pi d_P^2 \lambda f_{\lambda} \text{ ergs s}^{-1}$$





**Table 6.** Measurements on the broad lines of median spectra.

Object name (1)	$W(H\beta)^a$ (2)	$W(FeII\lambda 4570)^a$ (3)	$FWHM(FeII)^b$ (4)	$F(H\beta_{BC})/F(H\beta)^c$ (5)	$FWHM(H\beta_{BC})^d$ (6)	$\log M_{BH} H\beta_{BC}^e$ (7)	$\log L_{bol}/L_{Edd} H\beta_{BC}^e$ (8)
A1	72 <sup>+11</sup> <sub>-11</sub>	26 <sup>+3</sup> <sub>-4</sub>	2700 <sup>+1100</sup> <sub>-1100</sub>	1.00	...	...	...
A2	65 <sup>+10</sup> <sub>-52</sub>	49 <sup>+13</sup> <sub>-11</sub>	3700 <sup>+2000</sup> <sub>-1400</sub>	1.00	...	...	...
B1	86 <sup>+13</sup> <sub>-13</sub>	26 <sup>+5</sup> <sub>-6</sub>	5200 <sup>+2400</sup> <sub>-2300</sub>	0.27	4000	...	...
B2	70 <sup>+11</sup> <sub>-11</sub>	44 <sup>+8</sup> <sub>-14</sub>	5000 <sup>+800</sup> <sub>-1700</sub>	0.32	4000	...	...
A	61 <sup>+10</sup> <sub>-13</sub>	25 <sup>+7</sup> <sub>-7</sub>	2700 <sup>+1500</sup> <sub>-1200</sub>	1.00	...	...	...
M	67 <sup>+10</sup> <sub>-11</sub>	35 <sup>+7</sup> <sub>-6</sub>	3800 <sup>+1450</sup> <sub>-1200</sub>	1.00	...	...	...
MB	86 <sup>+10</sup> <sub>-13</sub>	31 <sup>+6</sup> <sub>-7</sub>	5000 <sup>+1600</sup> <sub>-1800</sub>	0.27	4100	...	...
43A	91 <sup>+10</sup> <sub>-20</sub>	36 <sup>+7</sup> <sub>-7</sub>	3000 <sup>+500</sup> <sub>-750</sub>	1.00	...	6.1	-0.74
44A	69 <sup>+15</sup> <sub>-15</sub>	38 <sup>+10</sup> <sub>-10</sub>	2600 <sup>+300</sup> <sub>-750</sub>	1.00	...	6.8	-0.47
45A	86 <sup>+10</sup> <sub>-20</sub>	43 <sup>+10</sup> <sub>-10</sub>	2800 <sup>+750</sup> <sub>-500</sub>	1.00	...	7.8	-0.43
46A	80 <sup>+10</sup> <sub>-10</sub>	47 <sup>+10</sup> <sub>-10</sub>	3000 <sup>+600</sup> <sub>-600</sub>	1.00	...	8.6	-0.26
47A	68 <sup>+10</sup> <sub>-11</sub>	30 <sup>+8</sup> <sub>-8</sub>	3000 <sup>+1400</sup> <sub>-1200</sub>	1.00	...	9.6	-0.20
48A	60 <sup>+11</sup> <sub>-11</sub>	27 <sup>+5</sup> <sub>-8</sub>	3800 <sup>+1500</sup> <sub>-1100</sub>	1.00	...	10.3	+0.11
43B	130 <sup>+20</sup> <sub>-20</sub>	8 <sup>+10</sup> <sub>-7</sub>	... <sup>f</sup>	0.59	4600	7.1	-0.68
44B	125 <sup>+10</sup> <sub>-30</sub>	38 <sup>+5</sup> <sub>-20</sub>	5600 <sup>+600</sup> <sub>-1800</sub>	0.49	4700	7.7	-1.37
45B	111 <sup>+15</sup> <sub>-20</sub>	29 <sup>+5</sup> <sub>-15</sub>	4900 <sup>+300</sup> <sub>-800</sub>	0.35	4400	8.4	-0.98
46B	93 <sup>+10</sup> <sub>-20</sub>	22 <sup>+5</sup> <sub>-10</sub>	5900 <sup>+350</sup> <sub>-1200</sub>	0.37	4800	9.1	-0.73
47B	92 <sup>+13</sup> <sub>-14</sub>	38 <sup>+7</sup> <sub>-7</sub>	4900 <sup>+1600</sup> <sub>-2000</sub>	0.27	4000	9.6	-0.24
48B	75 <sup>+9</sup> <sub>-11</sub>	12 <sup>+3</sup> <sub>-3</sub>	4600 <sup>+1200</sup> <sub>-1700</sub>	0.23	4300	10.3	+0.03

<sup>a</sup> Equivalent width of  $H\beta$  ( $H\beta_{BC} + H\beta_{VBC}$ ) and  $FeII\lambda 4570$  in  $\text{\AA}$   $\pm 2\sigma$  confidence level uncertainty. Note that those values have been computed on median spectra with flux normalized to unity at  $\lambda = 5100 \text{ \AA}$ . Considering that the continuum shape is not flat, but that there is however little dispersion in continuum shape across the median spectra, it is  $W(H\beta) \approx I(H\beta)/1.1$ , and  $W(FeII\lambda 4570) \approx I(FeII\lambda 4570)/1.25 \text{ \AA}$ . <sup>b</sup>  $FWHM$  of lines in the blend in units of  $\text{km s}^{-1}$  computed by `specfit` as for the individual sources. Uncertainty is at  $\pm 2\sigma$  confidence level. See text for details. <sup>c</sup> Intensity ratio of the  $H\beta_{BC}$  to total  $H\beta$  line emission i.e.,  $H\beta_{BC}$  and  $H\beta_{VBC}$ . <sup>d</sup>  $FWHM$  of the  $H\beta_{BC}$  component i.e., after removing  $H\beta_{VBC}$ . <sup>e</sup> Logarithm of  $M_{BH}$ , in solar masses, and of  $L_{bol}/L_{Edd}$ . Values have been computed following Paper II, using the  $FWHM(H\beta_{BC})$  reported in Col. (6), and assuming the average bin luminosity. Values are therefore only indicative. No  $M_{BH}$  or  $L_{bol}/L_{Edd}$  has been computed for median in spectral types since they are normalized median spectra made regardless of their luminosity. <sup>f</sup>  $FeII_{opt}$  too faint for  $FWHM$  to be meaningfully constrained.

